

$$\frac{\partial}{\partial t} f(v_1, t) = -\frac{2}{3} (\pi \sigma)^2 f(v_1, t) \int_0^\infty dv_2 v_2^2 f(v_2, t) \left[\frac{(v_1+v_2)^3 - |v_1-v_2|^3}{v_1 v_2} \right] \quad (1)$$

$$f(v_1, t) = \sum_{i=1}^N \frac{n_i(t)}{4\pi c_i^2} \delta(v_1 - c_i) \quad ; \quad c_1 > 0, \quad c_j > c_i \quad \forall j > i. \quad (2)$$

$$\Rightarrow \frac{1}{4\pi} \sum_{i=1}^N \frac{\dot{n}_i}{c_i^2} \delta(v_1 - c_i) = -\frac{2}{3} \frac{(\pi \sigma)^2}{4\pi} \sum_{i=1}^N \frac{n_i}{c_i^2} \delta(v_1 - c_i) \int_0^\infty dv_2 v_2^2 \left[\frac{(v_1+v_2)^3 - |v_1-v_2|^3}{v_1 v_2} \right] \frac{1}{4\pi} \sum_{i=1}^N \frac{n_i}{c_i^2} \delta(v_2 - c_i) \quad (3)$$

Coefficient de $\delta(v_1 - c_i)$:

$$\frac{1}{4\pi} \frac{\dot{n}_i}{c_i^2} = -\frac{2}{3} (\pi \sigma)^2 \frac{1}{4\pi} \frac{n_i}{c_i^2} \int_0^\infty dv_2 v_2^2 \left[\frac{(v_1+v_2)^3 - |v_1-v_2|^3}{v_1 v_2} \right] \frac{1}{4\pi} \sum_{j=1}^N \frac{n_j}{c_j^2} \delta(v_2 - c_j) \quad ; \quad |v_1| = c_i$$

$$\Rightarrow \dot{n}_i = -\frac{2}{3} \frac{(\pi \sigma)^2}{4\pi} n_i \sum_{j=1}^N \frac{n_j}{c_j^2} \int_0^\infty dv_2 v_2^2 \left[\frac{(v_1+v_2)^3 - |v_1-v_2|^3}{v_1 v_2} \right] \delta(v_2 - c_j) \Big|_{|v_1|=c_i}$$

$$= \int_0^\infty dv_2 v_2^2 \left[\frac{(c_i+v_2)^3 - |c_i-v_2|^3}{c_i v_2} \right] \delta(v_2 - c_j)$$

$$= c_j^2 \frac{(c_i+c_j)^3 - |c_i-c_j|^3}{c_i c_j}$$

$$= -\frac{2}{3} \frac{(\pi \sigma)^2}{4\pi} n_i \sum_{j=1}^N \frac{n_j}{c_j^2} \frac{c_j^2}{c_i c_j} \left((c_i+c_j)^3 - |c_i-c_j|^3 \right) \quad ; \quad \text{cas: } \begin{cases} i > j \\ i < j \\ i = j \end{cases}$$

$$= -\frac{\pi}{6} \sigma^2 n_i \sum_{j=1}^N n_j \frac{(c_i+c_j)^3 - |c_i-c_j|^3}{c_i c_j} \quad (3b)$$

$$= -\frac{\pi}{6} \sigma^2 n_i \left[\sum_{j=1}^{i-1} n_j \frac{(c_i+c_j)^3 - |c_i-c_j|^3}{c_i c_j} + \sum_{j=i}^i n_j \frac{(c_i+c_j)^3 - |c_i-c_j|^3}{c_i c_j} + \sum_{j=i+1}^N n_j \frac{(c_i+c_j)^3 - |c_i-c_j|^3}{c_i c_j} \right]$$

$i > j \Rightarrow c_i > c_j$
 Δ si $i=1$: si la somme n'est pas définie, alors on oublie ce terme.
 $c_i > c_j \Rightarrow |c_i - c_j| = c_i - c_j$
 $\Rightarrow \sum_{j=1}^{i-1} n_j \frac{(c_i+c_j)^3 - (c_i-c_j)^3}{c_i c_j}$
 $= \sum_{j=1}^{i-1} n_j \frac{1}{c_i c_j} (6c_i^2 c_j + 2c_j^3)$
 $= \sum_{j=1}^{i-1} n_j \left(6c_i \frac{c_j}{c_j} + 2 \frac{c_j^3}{c_i} \right) \frac{1}{c_j}$

$i = j = i \Rightarrow c_i = c_j$
 $= n_i \frac{(2c_i)^3}{c_i c_i}$
 $= 8 n_i c_i$

$i < j \Rightarrow c_i < c_j$
 Δ si $i=N$: " - ".
 $c_i < c_j \Rightarrow |c_i - c_j| = c_j - c_i$
 $= \sum_{j=i+1}^N n_j \frac{(c_i+c_j)^3 - (c_j-c_i)^3}{c_i c_j}$
 $= \sum_{j=i+1}^N n_j \frac{1}{c_i c_j} (2c_i^3 + 6c_i c_j^2)$
 $= \sum_{j=i+1}^N n_j \left(2 \frac{c_i^3}{c_j} + 6 c_j \frac{c_i^2}{c_j} \right) \frac{1}{c_j}$

$$= -\frac{\pi}{3} \sigma^2 n_i \left[\sum_{j=1}^{i-1} n_j \left(3c_i + \frac{c_j^2}{c_i} \right) + 4 n_i c_i + \sum_{j=i+1}^N n_j \left(\frac{c_i^2}{c_j} + 3c_j \right) \right] \cdot \frac{c_N}{c_N}$$

$$= -\frac{\pi}{3} \sigma^2 c_N n_i \left[\sum_{j=1}^{i-1} n_j \left(3 \frac{c_i}{c_N} + \frac{c_j^2}{c_i c_N} \right) + 4 n_i \frac{c_i}{c_N} + \sum_{j=i+1}^N n_j \left(\frac{c_i^2}{c_j c_N} + 3 \frac{c_j}{c_N} \right) \right] \cdot \frac{\pi \sigma^2}{3} c_N$$

$$\Rightarrow \dot{n}_i = -n_i \left[\sum_{j=1}^{i-1} n_j \left(3 \frac{c_i}{c_N} + \frac{c_j^2}{c_i c_N} \right) \frac{c_i}{c_N} + \sum_{j=i+1}^N n_j \left(\frac{c_i^2}{c_j c_N} + 3 \frac{c_j}{c_N} \right) \right] \quad ; \quad \delta_c = \frac{c_i}{c_N} < 1 \quad \forall i < N$$

$$\dot{n}_i = -n_i \left[\sum_{j=1}^{i-1} n_j \left(3 \delta_c + \frac{c_j}{c_i} \delta_j \right) + 4 n_i \delta_c + \sum_{j=i+1}^N n_j \left(3 \delta_j + \frac{c_i}{c_j} \delta_c \right) \right] \quad (4)$$

Vérification: cas particulier N=2:

$$\dot{n}_i = -n_i \left[\sum_{j=1}^{i-1} n_j (3\delta_{ij} + \frac{c_j}{c_i} \delta_j) + 4n_i \delta_i + \sum_{j=i+1}^2 n_j (3\delta_j + \frac{c_i}{c_j} \delta_i) \right]$$

$$\Rightarrow \begin{cases} \dot{n}_1 = -n_1 \left[\sum_{j=1}^0 + 4n_1 \delta_1 + \sum_{j=2}^2 n_j (3\delta_j + \frac{c_1}{c_j} \delta_1) \right] = -n_1 [4n_1 \delta_1 + n_2 (3 + \delta_1^2)] \\ \dot{n}_2 = -n_2 \left[\sum_{j=1}^1 n_j (3\delta_j + \frac{c_j}{c_2} \delta_j) + 4n_2 \delta_2 + \sum_{j=3}^2 \right] = -n_2 [n_1 (3 + \delta_1^2) + 4n_2] \end{cases}$$

n1=X
n2=Y
delta=delta

$$\begin{cases} \dot{X} = -4X^2\delta - XY(3+\delta^2) \\ \dot{Y} = -4Y^2 - XY(3+\delta^2) \end{cases}$$

Ok, ce sont les mêmes équations que précédemment établies.

Introduisons

$$\delta_{ij} = \frac{c_i}{c_j} \leq 1 \quad \forall j > i \quad \text{t.q.} \quad \delta_{iN} = \delta_i$$

et il faut formaliser les sommes: somme 1 indéfinie si i-1 < 1, i.e. si i < 2, donc il faut que i > 1.
Somme 2 indéfinie si N < i+1, i.e. si i > N-1, donc il faut que i < N.

$$\theta(x-a) = \begin{cases} 0 & \text{si } x \leq a \\ 1 & \text{si } x > a \end{cases}$$

$$\Rightarrow \dot{n}_i = -n_i \left[\theta(i-1) \sum_{j=1}^{i-1} n_j (3\delta_{ij} + \delta_{ij} \delta_j) + 4n_i \delta_i + \theta(N-i) \sum_{j=i+1}^N n_j (3\delta_j + \delta_{ij} \delta_i) \right] \quad (5)$$

ou bien de façon équivalente:

$$\begin{cases} \dot{n}_1 = -n_1 \left[4n_1 \delta_1 + \sum_{j=2}^N n_j (3\delta_j + \delta_{1j} \delta_1) \right] \\ \dot{n}_i = -n_i \left[\sum_{j=1}^{i-1} n_j (3\delta_{ij} + \delta_{ij} \delta_j) + 4n_i \delta_i + \sum_{j=i+1}^N n_j (3\delta_j + \delta_{ij} \delta_i) \right], \quad i=2, \dots, N-1 \\ \dot{n}_N = -n_N \left[\sum_{j=1}^{N-1} n_j (3 + \delta_j^2) + 4n_N \right] \end{cases} \quad (6)$$

Avec: $\delta_{ij} = \frac{c_i}{c_j} < 1 \quad \forall j > i; \quad c_i < c_{i+1} \quad \forall i=1, \dots, N-1$
 $\delta_i = \delta_{iN}; \quad \delta_i < \delta_{i+1} \quad \forall i=1, \dots, N-1$

Argument pour dire que n1(t) ~ 1/4delta1 tau : comme les vitesses des espèces plus rapides sont plus grandes que c1, alors ces espèces vont avoir tendance à décroître plus rapidement (on le voit des études 2d, 3d, par 2 c.i.), par conséquent les non linéarités du type ninj peuvent être négligées, ∀ ij ≠ 1. Si ni = O(ε) ∀ i ≠ 1, alors:

$$\begin{aligned} n_i n_j &= O(\epsilon^2) \quad \forall ij \neq 1 \\ n_i n_j &= O(\epsilon) \quad \forall j \neq 1 \\ n_i^2 &= O(1) \end{aligned}$$

⇒ on obtient par n1:

$$\dot{n}_1 = -4\delta_1 n_1^2 \quad ; \quad \frac{d}{dt} \left(\frac{1}{n_1} \right) = -\frac{1}{n_1^2} \dot{n}_1$$

$$\Rightarrow \frac{d}{dt} \left(\frac{1}{n_1} \right) = -4\delta_1$$

$$\Rightarrow \frac{n_1(\tau)}{n_1(0)} \approx \frac{1}{4\delta_1 \tau} \quad ; \quad \delta_1 = \frac{c_1}{c_N} \quad (7)$$

VÉRIFIÉ NUMÉRIQUEMENT

Généralisation en dimension $d \geq 2$

$$\begin{cases} \frac{\partial}{\partial t} f(v_i, t) = -\sigma^{d-1} \beta_d f(v_i, t) \int_{\mathbb{R}^d} dv_2 |v_{12}| f(v_2, t) & ; \beta_d = \int d\theta \theta (\sigma \hat{v}_{12}) (\sigma \hat{v}_{12}) = \frac{\pi^{\frac{d-1}{2}}}{\Gamma(\frac{d+1}{2})} \end{cases} \quad (8)$$

$$\begin{cases} f(v_i, t) = \sum_{i=1}^N \frac{n_i(t)}{S_d c_i^{d-1}} \delta(v - c_i) & ; S_d = \int_{\mathbb{R}^d} d\Omega = \frac{2\pi^{\frac{d}{2}}}{\Gamma(\frac{d}{2})} \end{cases} \quad (9)$$

(9) dans (8) \Rightarrow

$$\frac{1}{S_d} \sum_{i=1}^N \frac{\dot{n}_i}{c_i^{d-1}} \delta(v_1 - c_i) = -\sigma^{d-1} \beta_d \frac{1}{S_d} \sum_{i=1}^N \frac{n_i}{c_i^{d-1}} \delta(v_1 - c_i) \int_{\mathbb{R}^d} dv_2 |v_{12}| \frac{1}{S_d} \sum_{j=1}^N \frac{n_j}{c_j^{d-1}} \delta(v_2 - c_j) \quad (10)$$

Coefficients de $\delta(v_1 - c_i)$:

$$\frac{1}{S_d} \frac{\dot{n}_i}{c_i^{d-1}} = -\sigma^{d-1} \beta_d \frac{1}{S_d} \frac{n_i}{c_i^{d-1}} \int_{\mathbb{R}^d} dv_2 |v_{12}| \frac{1}{S_d} \sum_{j=1}^N \frac{n_j}{c_j^{d-1}} \delta(v_2 - c_j) \Big|_{|v_1|=c_i}$$

$$\Rightarrow \dot{n}_i = -\frac{\sigma^{d-1} \beta_d}{S_d} n_i \sum_{j=1}^N \frac{n_j}{c_j^{d-1}} \int_{\mathbb{R}^d} dv_2 |v_{12}| \delta(v_2 - c_j) \Big|_{|v_1|=c_i} =: I_{ij}$$

$$\int_{\mathbb{R}^d} dv \sqrt{v^2 + c^2 - 2vc \cos \theta} \delta(v - c) = \int_0^\pi d\theta \int_0^\infty dv \sqrt{v^2 + c^2 - 2vc \cos \theta} (nc)^{d-2} = \int_0^\pi d\theta \int_0^\infty dv \sqrt{c^2 + c^2 - 2cc \cos \theta} (nc)^{d-2} = \int_0^\pi d\theta \int_0^\infty dv \sqrt{2c^2(1 - \cos \theta)} (nc)^{d-2}$$

Avec I_{ij} déjà calculé dans généralisation à d dimensions du cas à deux vitesses $N=2$

$$I_{ij} = 2\pi J_d \sqrt{c_i^2 + c_j^2} c_i^{d-1} F_d(c_i, c_j), \quad \forall i, j = 1, \dots, N$$

$$; J_d = \begin{cases} 1/\pi, & \text{si } d=2 \\ 1, & \text{si } d=3 \\ \frac{d-3}{\Gamma(\frac{d-1}{2})} \int_0^\pi d\theta \sin^k(\theta) & \text{si } d > 3 \end{cases} \quad (11)$$

ou bien: $F_d(c_i, c_j) = \int_0^\pi d\theta \sqrt{1 - 2 \frac{c_i c_j}{c_i^2 + c_j^2} \cos \theta} (\sin \theta)^{d-2}$

$i=j \Rightarrow F_d = \int_0^\pi d\theta \sqrt{1 - \cos \theta} (\sin \theta)^{d-2} = \int_0^\pi d\theta \sqrt{2(1 - \cos \theta)} (\sin \theta)^{d-2} = \int_0^\pi d\theta \sqrt{4 \sin^2(\frac{\theta}{2})} (\sin \theta)^{d-2} = \int_0^\pi d\theta 2 \sin(\frac{\theta}{2}) (\sin \theta)^{d-2}$

$$F_d(c_i, c_j) = \int_{-1}^1 dx \sqrt{1 - 2 \frac{c_i c_j}{c_i^2 + c_j^2} x} / (1-x^2)^{\frac{d-3}{2}} = F_d(c_j, c_i)$$

$$F_d = F_d(c_i, c_i) = \int_{-1}^1 dx \sqrt{1-x} / (1-x^2)^{\frac{d-3}{2}}$$

Si on fait le changement d'échelle de temps (défini de sorte que $n_i \sim 1/\sqrt{\tau}$)

$$\tau = t \sigma^{d-1} \frac{\beta_d J_d F_d \pi}{S_d \sqrt{2}} c_N \quad (12)$$

alors on retrouve bien celui en trois dimensions en posant $d=3$ dans (12):

$$\tau(d=3) = t \sigma^2 \frac{\pi \cdot 1}{4\pi} \frac{4}{3} \frac{\pi}{\sqrt{2}} c_N = t \frac{\pi \sigma^2}{3} c_N$$

(11) dans (10b) \Rightarrow

$$\dot{n}_i = -\frac{\sigma^{d-1} \beta_d}{S_d} n_i \sum_{j=1}^N \frac{n_j}{c_j^{d-1}} 2\pi J_d \sqrt{c_i^2 + c_j^2} c_j^{d-1} F_d(c_i, c_j) \quad \left| \cdot \frac{c_N}{c_N} \cdot \frac{F_d}{2\sqrt{2}} \cdot \frac{2\sqrt{2}}{F_d} \right.$$

$$= -\sigma^{d-1} \beta_d J_d \frac{F_d \pi}{S_d} c_N \frac{2\sqrt{2}}{F_d} \frac{1}{c_N} n_i \sum_{j=1}^N n_j \sqrt{c_i^2 + c_j^2} F_d(c_i, c_j)$$

$$\Rightarrow \dot{n}_i = -\frac{2\sqrt{2}}{F_d} n_i \sum_{j=1}^N n_j \sqrt{\frac{c_i^2}{c_N^2} + \frac{c_j^2}{c_N^2}} F_d(c_i, c_j)$$

$$\Rightarrow \dot{n}_i = -\frac{2\sqrt{2}}{F_d} n_i \sum_{j=1}^N n_j \sqrt{c_i^2 + c_j^2} F_d(c_i, c_j) \quad (13)$$

Argument sur la décroissance $n_i \sim 1/\sqrt{\tau}$ en dimension $d \geq 2$:

$$\dot{n}_i \approx -2\sqrt{2} n_i \sum_{j=1}^N n_j \sqrt{c_i^2 + c_j^2} \frac{F_d(c_i, c_j)}{F_d} \delta_{j,1}$$

$$= -4\sqrt{2} n_i^2$$

$$\Rightarrow n_i \approx \frac{1}{4\sqrt{2} \tau}$$

et aussi: changement de variable inverse:
 $n_i = \frac{B}{c_N^d} \sim \frac{B}{c_N^d} c_N \Rightarrow n_i(t) \sim \frac{B}{A c_N t} c_N \sim \frac{B}{A c_N t}$
 $\tau = A c_N t$
 indép de toutes les vitesses sauf c_i, c_j

Question: est-ce que (13) est bien équivalent à (6) (par d'erreur de calcul)? Relation (13) pour $d=3$ donne: (avec $F_3 = 4/3 \sqrt{2}$)

$$\dot{n}_i = -\frac{2\sqrt{2}}{4/3 \sqrt{2}} n_i \sum_{j=1}^N n_j \sqrt{c_i^2 + c_j^2} \int_{-1}^1 dx \sqrt{1 - 2 \frac{c_i c_j}{c_i^2 + c_j^2} x}$$

$$= \frac{1}{c_N} \int_{-1}^1 dx \sqrt{c_i^2 + c_j^2 - 2c_i c_j x}$$

$$= \frac{1}{c_N} \frac{1}{3} \frac{(c_i + c_j)^2 - |c_i - c_j|^3}{c_i c_j}$$

$$\Rightarrow \dot{n}_i = -\frac{\gamma}{2} n_i \sum_{j=1}^N \frac{1}{\gamma} \frac{1}{c_N} \frac{(c_i+c_j)^3 - |c_i-c_j|^3}{c_i c_j} ; \gamma = \pm \frac{\pi \sigma^2}{3} c_N$$

$$= -\frac{1}{2} n_i \frac{1}{\cancel{\gamma}} \frac{\pi \sigma^2}{3} \sum_{j=1}^N n_j \frac{(c_i+c_j)^3 - |c_i-c_j|^3}{c_i c_j}$$

$$= -\frac{\pi}{6} \sigma^2 n_i \sum_{j=1}^N n_j \frac{(c_i+c_j)^3 - |c_i-c_j|^3}{c_i c_j}$$

Ce qui est bien la même relation que (36), donc (6) s'obtient bien de (13).

Remarques: - structure de l'équation générale (13):

$$\dot{\underline{n}} = - \underline{c} \underline{n} \underline{A} \underline{n} ; \underline{n} = (n_1, \dots, n_N) ; \underline{A} = \{a_{ij}\}_{i,j=1}^N ; a_{ij} = \frac{2\sqrt{2}}{\sqrt{3}} \sqrt{|c_i^2 + c_j^2|} F_d(c_i, c_j) > 0 \forall i, j$$

- études à faire:

- preuve analytique de la décroissance algébrique des n_i , $c > 1$
- cas à 1 vitesse $c_1 = 0$

Essayons de prouver analytiquement la décroissance en loi de puissance : on a une équation du type :

$$\dot{n}_i = -n_i \sum_{j=1}^N a_{ij} n_j \quad ; \quad a_{ij} = \frac{3}{2} \sqrt{\delta_i^{-1} \delta_j^{-2}} F_0(c_i, c_j) > 0 \quad (1)$$

$$\Rightarrow \boxed{\dot{n} = - \langle n | A n \rangle} \quad ; \quad A = \{a_{ij}\}_{i,j=1}^N$$

par analogie au cas $\dot{n} = -an^2 \Rightarrow n \sim 1/akt$, est-ce qu'il n'existe pas mthm. similaires ?

(2)

Supposons que pour les temps longs :

$$\begin{cases} n_i \approx B_i t^{-\beta_i} & ; \beta_i > \beta_j \quad \forall i > j & ; \beta_1 = 1 ; B_i > 0 \quad \forall i \\ \dot{n}_i \approx -\beta_i \frac{n_i}{t} = -\beta_i B_i t^{-\beta_i-1} & & \text{(dans } \beta_i > 1 \forall i > 1) \end{cases}$$

$$\stackrel{(1)}{\Rightarrow} -\beta_i \frac{n_i}{t} = -n_i \sum_{j=1}^N a_{ij} n_j$$

$$\Rightarrow -\beta_i B_i t^{-\beta_i-1} = -\beta_i t^{-\beta_i} \sum_{j=1}^N a_{ij} B_j t^{-\beta_j}$$

$$\Rightarrow \beta_i t^{-\beta_i-1} = t^{-\beta_i} \sum_{j=1}^N a_{ij} B_j t^{-\beta_j}$$

$$\Rightarrow \beta_i t^{-1} = \sum_{j=1}^N a_{ij} B_j t^{-\beta_j}$$

$$\Rightarrow \boxed{\beta_i = \sum_{j=1}^N a_{ij} B_j t^{1-\beta_j}}$$

$$; \quad 1-\beta_j < 0 \quad \forall j > 1 ; \quad 1-\beta_j = 0 \quad \text{si } j=1.$$

~~Volterra
Lotka-généralisé
Baughoffer~~

~~Motbauer, Sigmund
Evolutionary Games and Population Dynamics
Cambr. Univ. Press. 1988~~

~~→ E.W. Montroll, Rev. Mod. Phys., 1965~~

~~PROLA~~

Il faut voir, il existe une telle solution avec $a_{ij} > 0, B_j > 0$ et $\beta_1 = 1, \beta_j > 1 \quad \forall j > 1$. Supposons que dans le régime $t \rightarrow \infty$ on puisse négliger tout ce qui décroît plus vite que le terme i , i.e. on suppose $\beta_i > \beta_j \quad \forall i > j$, alors :

$$\beta_i = \sum_{j=1}^i a_{ij} B_j t^{1-\beta_j}$$

avec :

$$\beta_1 = 1$$

$$\Rightarrow \beta_2 = a_{21} B_1 + \underbrace{a_{22} B_2 t^{1-\beta_2}}_{\rightarrow 0 \text{ n'a été l'hyp. } \beta_2 > 1} \Rightarrow \beta_2 = a_{21} B_1 > 0 ; \beta_2 > 1 ? \text{ hyp. auto consistante?}$$

$$\Rightarrow \beta_3 = a_{31} B_1 + \underbrace{a_{32} B_2 t^{1-\beta_2}}_0 + \underbrace{a_{33} B_3 t^{1-\beta_3}}_0$$

$$\Rightarrow \beta_1 = 1 ; \beta_i = \beta_2 \quad \forall i \geq 2 \Rightarrow \Leftarrow \quad \Leftarrow$$

Le même argument que celui utilisé pour $N=2$ ne semble pas marcher ici...

Question: pourquoi θ_{d-2} et non θ_1 dans l'intégration?

Réponse: par construction des coordonnées sphériques dans \mathbb{R}^d

Explications: soient $\underline{v}_1, \underline{v}_2 \in \mathbb{R}^d$, soit l'intégration: (on suppose $|\underline{v}_1| = 1$ pour alléger l'écriture)

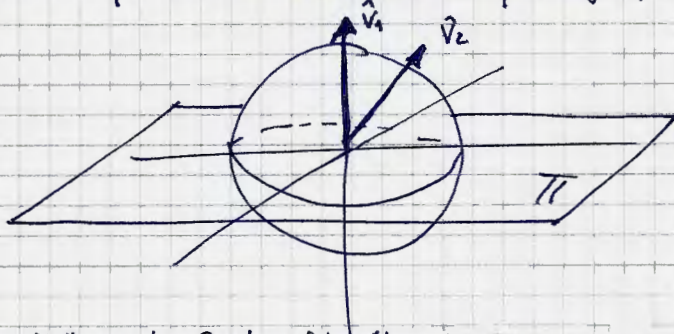
$$I = \int_{\mathbb{R}^d} d\underline{v}_2 \delta(|\underline{v}_2|-1) \underline{v}_1 \cdot \underline{v}_2 = \int_0^{2\pi} d\varphi \int_0^\pi d\theta_1 \dots \int_0^\pi d\theta_{d-1} \underline{v}_1 \cdot \underline{v}_2 \prod_{k=1}^{d-2} (\sin \theta_k)^{k-1}$$

$$= 2\pi \int_0^\pi d\theta_1 \dots \int_0^\pi d\theta_{d-1} \underbrace{|\underline{v}_1 \cdot \underline{v}_2|}_{=1} \cos \theta_{d-2} \prod_{k=1}^{d-2} (\sin \theta_k)^{k-1}$$

pourquoi est-ce l'angle θ_{d-2} qui intervient ici? Réponse: construisons les coordonnées sphériques dans \mathbb{R}^d , commençons par x_1 .

$$x_1 = \dots ? \dots = f(\theta_1, \dots, \theta_{d-2}, \varphi) \quad \text{trouver } f.$$

Il faut commencer par chercher une direction privilégiée, prenons \underline{v}_1 :



(illustration dans \mathbb{R}^3)

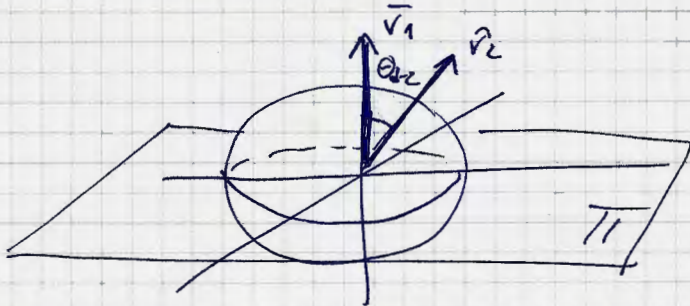
$$\Pi \perp \text{vect}\{\underline{v}_1\}$$

\downarrow
= ensemble des combinaisons linéaires.

L'idée est essentiellement Gram-Schmidt.

On projette les composantes de $\underline{v}_2 \in \mathbb{R}^d$ sur la direction \underline{v}_1 , i.e. dans un espace de dimension 1 $\text{vect}\{\underline{v}_1\}$. Pour ceci, on doit introduire l'angle θ_{d-2} entre \underline{v}_1 et \underline{v}_2 qui est donc défini comme l'angle entre

\underline{v}_1 et \underline{v}_2 dans le sous-espace $\text{vect}\{\underline{v}_1, \underline{v}_2\}$. Cet angle θ_{d-2} existe donc toujours.



$$x_1 = \cos \theta_{d-2}$$

Par x_2 , on doit construire l'orthogonal à $\text{vect}\{\underline{v}_1\}$ qui lui soit complémentaire dans \mathbb{R}^d , i.e.

$$\text{vect}\{\underline{v}_1\} \oplus \underbrace{\text{vect}\{\underline{v}_1\}^\perp}_{=\pi \text{ dans le demi-espace}} = \mathbb{R}^d \quad \Rightarrow \quad x_2 \sim \sin \theta_{d-2}$$

A l'intérieur de l'orthogonal $\text{vect}\{\underline{v}_1\}^\perp$, on en extrait un sous-espace de dimension 1 grâce à θ_{d-3} :

$$x_2 \sim \sin \theta_{d-2} \cdot \cos \theta_{d-3}$$

et son orthogonal de dimension 1

$$x_3 \sim \sin \theta_{d-2} \sin \theta_{d-3} \cos \theta_{d-4}$$

et ainsi de suite jusqu'au dernier angle θ :

$$x_{d-1} \sim \sin \theta_{d-2} \dots \sin \theta_1 \cos \theta$$

$$x_d \sim \sin \theta_{d-2} \dots \sin \theta_1 \sin \theta$$

On a ainsi décomposé \mathbb{R}^d en somme directe de sous-espaces orthogonaux. On voit donc que l'angle entre \underline{v}_1 et \underline{v}_2 doit être θ_{d-2} car cet angle définit la projection de $\underline{v}_2 \in \mathbb{R}^d$ dans le sous-espace de dimension 1 engendré par \underline{v}_1 . Prendre un autre angle que θ_{d-2} reviendrait à faire

n'importe quoi (même par une projection de n'importe quel sous-espace car il n'y aurait qu'un angle au lieu de plusieurs...)

Dimension arbitraire $d \geq 3$: cas général $0 < c_1 < c_2$ (+ convention $c_1 = 0 \text{ à } 1 \text{ m}$)

$$\frac{\partial}{\partial t} f(v,t) = -\sigma^{d-1} \beta_d f(v,t) \int_{\mathbb{R}^d} dv_2 |v_1| f(v_2,t) \quad ; \beta_d = \int d\vec{\theta} \Theta(\vec{\sigma} \cdot \vec{v}_2) (\vec{\sigma} \cdot \vec{v}_2) = \pi^{\frac{d-1}{2}} \Gamma\left(\frac{d+1}{2}\right)$$

Preliminaire :

$$I_{ij} = \int_{\mathbb{R}^d} dv_2 |v_2| \delta(v_2 - c_i) \stackrel{|v_1|=c_j}{=} \int_{\mathbb{R}^d} dv_2 |\hat{v}_1 c_j - \hat{v}_2 c_i| \delta(v_2 - c_i) \quad ; w = v_2/c_i \quad ; dv_2 = c_i^d dw$$

$$= \int_{\mathbb{R}^d} dw c_i^d |\hat{v}_1 c_j - \hat{w} c_i| \delta\left(\frac{w c_i}{c_i} - 1\right) \stackrel{= \frac{1}{c_i} \delta(w-1)}{=} c_i^{d-1} \int_{\mathbb{R}^d} dw |\hat{v}_1 c_j - \hat{w} c_i| \delta(w-1)$$

1) $i=j$: $I_{ii} = c_i^{d-1} \int_{\mathbb{R}^d} dw c_i |\hat{v}_1 - \hat{w}| \delta(w-1)$

$$= c_i^d \int_{\mathbb{R}^d} dw v_2 \sqrt{1 - \cos \theta_{d-2}} \delta(w-1)$$

$$= \sqrt{2} c_i^d \int_0^{2\pi} d\varphi \int_0^\pi d\theta_1 \dots \int_0^\pi d\theta_{d-2} \sqrt{1 - \cos \theta_{d-2}} \prod_{k=1}^{d-2} (\sin \theta_k)^k = (\sin \theta_k)^k$$

$$= \sqrt{2} c_i^d 2\pi \underbrace{\left(\prod_{k=1}^{d-3} \int_0^\pi d\theta_k (\sin \theta_k)^k \right)}_{:= J_d} \int_0^\pi d\theta_{d-2} \sqrt{1 - \cos \theta_{d-2}} (\sin \theta_{d-2})^{d-2}$$

$x = \cos \theta_{d-2}$
 $dx = -\sin \theta_{d-2} d\theta$

$$= \int_{-1}^1 \frac{dx}{-\sin \theta_{d-2}} (\sin \theta_{d-2})^{d-2} \sqrt{1-x}$$

$\theta_{d-2} = \arccos(x)$
 $\sin(\arccos) = \frac{1}{\sqrt{1-x^2}}$

$$= \int_{-1}^1 dx (\sin \theta_{d-2})^{d-3} \sqrt{1-x}$$

$$= \int_{-1}^1 dx \sqrt{1-x} \left(\frac{1}{\sqrt{1-x^2}} \right)^{d-3}$$

$$= \int_{-1}^1 dx \frac{\sqrt{1-x}}{(1-x^2)^{\frac{d-3}{2}}}$$

$$= 2\pi \sqrt{2} c_i^d F_d J_d$$

$= F_d$: l'intégrale se calcule en fonction de fonctions hypergéométriques
Sauf : $d=5$. Non bon, je préfère l'expression générale F_d .

2) $i \neq j$: $I_{ij} = c_i^{d-1} \int_{\mathbb{R}^d} dw |\hat{v}_1 c_j - \hat{w} c_i| \delta(w-1)$

$$= c_i^{d-1} \int_{\mathbb{R}^d} dw \sqrt{c_i^2 + c_j^2 - 2c_i c_j \cos \theta_{d-2}} \delta(w-1)$$

$$= c_i^{d-1} \int_0^{2\pi} d\varphi \int_0^\pi d\theta_1 \dots \int_0^\pi d\theta_{d-2} \sqrt{c_i^2 + c_j^2 - 2c_i c_j \cos \theta_{d-2}} \prod_{k=1}^{d-2} (\sin \theta_k)^k$$

$$= c_i^{d-1} 2\pi \underbrace{\left(\prod_{k=1}^{d-3} \int_0^\pi d\theta_k (\sin \theta_k)^k \right)}_{:= J_d} \int_0^\pi d\theta_{d-2} \sqrt{c_i^2 + c_j^2 - 2c_i c_j \cos \theta_{d-2}} (\sin \theta_{d-2})^{d-2}$$

$x = \cos \theta_{d-2}$
 $dx = -\sin \theta_{d-2} d\theta$

$$= \int_{-1}^1 \frac{dx}{-\sin \theta_{d-2}} \sqrt{c_i^2 + c_j^2 - 2c_i c_j x} (\sin \theta_{d-2})^{d-2}$$

$$= \int_{-1}^1 dx (\sin \theta_{d-2})^{d-3} \sqrt{c_i^2 + c_j^2 - 2c_i c_j x}$$

$$= \int_{-1}^1 dx \frac{1}{(\sqrt{1-x^2})^{d-3}} \sqrt{c_i^2 + c_j^2} \sqrt{1 - 2 \frac{c_i c_j}{c_i^2 + c_j^2} x}$$

on bien tenir la forme en $\cos \theta_{d-2}$

$$= \sqrt{c_i^2 + c_j^2} \int_{-1}^1 dx \frac{\sqrt{1 - 2 \frac{c_i c_j}{c_i^2 + c_j^2} x}}{(1-x^2)^{\frac{d-3}{2}}}$$

$$= F_d(c_i, c_j)$$

avec : $F_d = F_d(1, 1)$

$$= c_i^{d-1} 2\pi J_d \sqrt{c_i^2 + c_j^2} F_d(c_i, c_j)$$

En résumé: $I_{ii} = 2\pi \int_{\mathbb{R}^d} C_i^{d-1} \sqrt{2} F_d(1,1)$
 $I_{ij} = 2\pi \int_{\mathbb{R}^d} C_i^{d-1} \sqrt{C_i^2 + C_j^2} F_d(C_i, C_j)$; $\lim_{C_i \rightarrow 0} F_d(C_i, C_j) = \int_{-1}^1 dx \frac{1}{(1-x^2)^{\frac{d-1}{2}}}$: diverge: ne peut pas prendre la limite à posteriori (2)

En fait $\lim_{C_i \rightarrow C_j} F_d(C_i, C_j) = F_d(C_i, C_i) = F_d$ est indép. de C_i ,
 donc on va noter: $\lim_{C_i \rightarrow C_j} F_d(C_i, C_j) = F_d(1,1) = F_d(C_i, C_i)$

$$\begin{cases} I_{ii} = 2\pi \int_{\mathbb{R}^d} C_i^{d-1} \sqrt{2} F_d \\ I_{ij} = 2\pi \int_{\mathbb{R}^d} C_i^{d-1} \sqrt{C_i^2 + C_j^2} F_d(C_i, C_j) \end{cases} \text{ t.q. } F_d(C_i, C_i) = F_d$$

Fonction de distribution:

$$f(v, t) = X \frac{1}{\int_{\mathbb{R}^d} C_i^{d-1}} \delta(v - c_1) + Y \frac{1}{\int_{\mathbb{R}^d} C_2^{d-1}} \delta(v - c_2) ; \int_{\mathbb{R}^d} d\Omega = \frac{2\pi^{d/2}}{\Gamma(d/2)}$$

Equation:

$$\dot{X} \frac{1}{\int_{\mathbb{R}^d} C_i^{d-1}} \delta(v - c_1) + \dot{Y} \frac{1}{\int_{\mathbb{R}^d} C_2^{d-1}} \delta(v - c_2) = -\sigma^{d-1} \beta_d \left[\frac{X}{\int_{\mathbb{R}^d} C_i^{d-1}} \delta(v - c_1) + \frac{Y}{\int_{\mathbb{R}^d} C_2^{d-1}} \delta(v - c_2) \right] \int_{\mathbb{R}^d} dv_2 |v_2| \left[\frac{X}{\int_{\mathbb{R}^d} C_i^{d-1}} \delta(v_2 - c_1) + \frac{Y}{\int_{\mathbb{R}^d} C_2^{d-1}} \delta(v_2 - c_2) \right]$$

Coefficients \Rightarrow equation par X et Y :

$$\begin{cases} \dot{X} \frac{1}{\int_{\mathbb{R}^d} C_i^{d-1}} = -\sigma^{d-1} \beta_d \frac{X}{\int_{\mathbb{R}^d} C_i^{d-1}} \int_{\mathbb{R}^d} dv_2 |v_2| \left[\frac{X}{\int_{\mathbb{R}^d} C_i^{d-1}} \delta(v_2 - c_1) + \frac{Y}{\int_{\mathbb{R}^d} C_2^{d-1}} \delta(v_2 - c_2) \right] \\ \dot{Y} \frac{1}{\int_{\mathbb{R}^d} C_2^{d-1}} = -\sigma^{d-1} \beta_d \frac{Y}{\int_{\mathbb{R}^d} C_2^{d-1}} \int_{\mathbb{R}^d} dv_2 |v_2| \left[\frac{X}{\int_{\mathbb{R}^d} C_i^{d-1}} \delta(v_2 - c_1) + \frac{Y}{\int_{\mathbb{R}^d} C_2^{d-1}} \delta(v_2 - c_2) \right] \end{cases}$$

$$\Rightarrow \begin{cases} \dot{X} = -\sigma^{d-1} \beta_d \left[\frac{X^2}{\int_{\mathbb{R}^d} C_i^{d-1}} \int_{\mathbb{R}^d} dv_2 |v_2| \delta(v_2 - c_1) + \frac{XY}{\int_{\mathbb{R}^d} C_2^{d-1}} \int_{\mathbb{R}^d} dv_2 |v_2| \delta(v_2 - c_2) \right] \\ \dot{Y} = -\sigma^{d-1} \beta_d \left[\frac{Y^2}{\int_{\mathbb{R}^d} C_2^{d-1}} \int_{\mathbb{R}^d} dv_2 |v_2| \delta(v_2 - c_2) + \frac{XY}{\int_{\mathbb{R}^d} C_i^{d-1}} \int_{\mathbb{R}^d} dv_2 |v_2| \delta(v_2 - c_1) \right] \end{cases}$$

$$I_{11} = 2\pi \int_{\mathbb{R}^d} C_i^{d-1} \sqrt{2} F_d ; I_{22} = 2\pi \int_{\mathbb{R}^d} C_2^{d-1} \sqrt{C_1^2 + C_2^2} F_d(C_1, C_2)$$

$$I_{22} = 2\pi \int_{\mathbb{R}^d} C_2^{d-1} \sqrt{2} F_d ; I_{12} = 2\pi \int_{\mathbb{R}^d} C_i^{d-1} \sqrt{C_1^2 + C_2^2} F_d(C_i, C_j)$$

$$\Rightarrow \begin{cases} \dot{X} = -\sigma^{d-1} \beta_d 2\pi \int_{\mathbb{R}^d} \left[X^2 \frac{C_i^{d-1} \sqrt{2} F_d}{\int_{\mathbb{R}^d} C_i^{d-1}} + XY \frac{C_2^{d-1} \sqrt{C_1^2 + C_2^2} F_d(C_1, C_2)}{\int_{\mathbb{R}^d} C_2^{d-1}} \right] \\ \dot{Y} = -\sigma^{d-1} \beta_d 2\pi \int_{\mathbb{R}^d} \left[Y^2 \frac{C_2^{d-1} \sqrt{2} F_d}{\int_{\mathbb{R}^d} C_2^{d-1}} + XY \frac{C_i^{d-1} \sqrt{C_1^2 + C_2^2} F_d(C_1, C_2)}{\int_{\mathbb{R}^d} C_i^{d-1}} \right] \end{cases}$$

$$\Rightarrow \begin{cases} \dot{X} = -\sigma^{d-1} \beta_d \frac{2\pi \int_{\mathbb{R}^d}}{\int_{\mathbb{R}^d}} \left[X^2 C_i \sqrt{2} F_d + XY \frac{\sqrt{C_1^2 + C_2^2} F_d(C_1, C_2)}{\sqrt{2}} \right] \cdot \frac{C_2}{C_1} \\ \dot{Y} = -\sigma^{d-1} \beta_d \frac{2\pi \int_{\mathbb{R}^d}}{\int_{\mathbb{R}^d}} \left[Y^2 C_2 \sqrt{2} F_d + XY \frac{\sqrt{C_1^2 + C_2^2} F_d(C_1, C_2)}{\sqrt{2}} \right] \cdot \frac{C_1}{C_2} \end{cases}$$

$$\Rightarrow \begin{cases} \dot{X} = -\sigma^{d-1} 2\pi \frac{\beta_d \int_{\mathbb{R}^d} F_d}{\int_{\mathbb{R}^d}} \sqrt{2} C_1 \left[X^2 \frac{C_1}{C_2} + XY \frac{\sqrt{C_1^2 + C_2^2}}{\sqrt{2} C_2} \frac{F_d(C_1, C_2)}{F_d} \right] \\ \dot{Y} = -\sigma^{d-1} 2\pi \frac{\beta_d \int_{\mathbb{R}^d} F_d}{\int_{\mathbb{R}^d}} \sqrt{2} C_2 \left[Y^2 + XY \frac{\sqrt{C_1^2 + C_2^2}}{\sqrt{2} C_2} \frac{F_d(C_1, C_2)}{F_d} \right] \end{cases} ; \tau = \sigma^{d-1} 2\pi \frac{\beta_d \int_{\mathbb{R}^d} F_d}{\int_{\mathbb{R}^d}} \sqrt{2} C_2$$

$$\Rightarrow \begin{cases} \dot{X} = -\delta X^2 - XY F \\ \dot{Y} = -Y^2 - XY F \end{cases}$$

$$\tau = \sigma^{d-1} 2\pi \frac{\beta_d \int_{\mathbb{R}^d} F_d}{\int_{\mathbb{R}^d}} \sqrt{2} C_2$$

$$\beta_d = \pi^{\frac{d-1}{2}} / \Gamma(\frac{d+1}{2}) ; \int_{\mathbb{R}^d} = \int_{\theta=0}^{2\pi} \int_{\phi=0}^{\pi} d\Omega (\sin\theta)^{d-2} ; \int_0^1 dx (1-x^2)^k = \pi \frac{\Gamma(k+1)}{\Gamma(1+k/2)}$$

$$\int_{\mathbb{R}^d} = 2\pi^{d/2} / \Gamma(d/2) ; F_d = \lim_{C_i \rightarrow C_j} F_d(C_i, C_j)$$

$$F_d(C_i, C_j) = \int_{-1}^1 dx \sqrt{1 - 2 \frac{C_i C_j}{C_i^2 + C_j^2} x} / (1-x^2)^{\frac{d-3}{2}}$$

$$F = \frac{\sqrt{C_1^2 + C_2^2}}{\sqrt{2} C_2} F_d(C_1, C_2) / F_d$$

Il faudrait vérifier avec le cas $C_1 = 0$...

Vérification avec le cas $C_1 = 0$:

$$f(v,t) = X \delta(v) + \frac{Y}{\int_{\mathbb{R}^d} \delta(v-c_2)} \delta(v-c_2)$$

$$\Rightarrow \begin{cases} \dot{X} = -\sigma^{d-1} \beta_d X \int_{\mathbb{R}^d} dv_2 |v_2| \left[X \delta(v_2) + \frac{Y}{\int_{\mathbb{R}^d} \delta(v_2-c_2)} \delta(v_2-c_2) \right] \\ \dot{Y} = -\sigma^{d-1} \beta_d \frac{Y}{\int_{\mathbb{R}^d} \delta(v_2-c_2)} \int_{\mathbb{R}^d} dv_2 |v_2| \left[X \delta(v_2) + \frac{Y}{\int_{\mathbb{R}^d} \delta(v_2-c_2)} \delta(v_2-c_2) \right] \end{cases}$$

$$\Rightarrow \begin{cases} \dot{X} = -\sigma^{d-1} \beta_d \left[X^2 \int_{\mathbb{R}^d} dv_2 |v_2| \delta(v_2) + \frac{XY}{\int_{\mathbb{R}^d} \delta(v_2-c_2)} \int_{\mathbb{R}^d} dv_2 |v_2| \delta(v_2-c_2) \right] & : f(v_1) \\ \dot{Y} = -\sigma^{d-1} \beta_d \left[\frac{Y^2}{\int_{\mathbb{R}^d} \delta(v_2-c_2)} \int_{\mathbb{R}^d} dv_2 |v_2| \delta(v_2-c_2) + XY \int_{\mathbb{R}^d} dv_2 |v_2| \delta(v_2) \right] & : f(v_2-c_2) \end{cases}$$

avec:

$$\int_{\mathbb{R}^d} \delta(v_1) \int_{\mathbb{R}^d} dv_2 |v_2| \delta(v_2) = \int_{\mathbb{R}^d} dv_2 |v_2| \delta(v_2) = 0$$

$$\begin{aligned} \int_{\mathbb{R}^d} \delta(v_1) \int_{\mathbb{R}^d} dv_2 |v_1-v_2| \delta(v_2-c_2) &= \int_{\mathbb{R}^d} dv_2 |v_2| \delta(v_2-c_2) & ; \underline{w} = v_2/c_2 ; dw = \frac{1}{c_2^d} dv_2 \\ &= c_2^d \int_{\mathbb{R}^d} dw |c_2 \hat{w}| \delta(c_2 w - c_2) \\ &= c_2^d \int_{\mathbb{R}^d} d\Omega \\ &= c_2^d S_d \end{aligned}$$

$$\int_{\mathbb{R}^d} \delta(v_1-c_2) \int_{\mathbb{R}^d} dv_2 |v_1-v_2| \delta(v_2) = \int_{\mathbb{R}^d} dv_2 |\hat{v}_1 c_2 - v_2| \delta(v_2) = |\hat{v}_1 c_2| = c_2$$

$$\begin{aligned} \int_{\mathbb{R}^d} \delta(v_1-c_2) \int_{\mathbb{R}^d} dv_2 |v_1-v_2| \delta(v_2-c_2) &= \int_{\mathbb{R}^d} dv_2 |\hat{v}_1 c_2 - v_2| \delta(v_2-c_2) ; \underline{w} = v_2/c_2 ; dw = \frac{1}{c_2^d} dv_2 \\ &= \int_{\mathbb{R}^d} dw c_2^d |\hat{v}_1 c_2 - \hat{w} c_2| \delta(w c_2 - c_2) \\ &= c_2^d \int_{\mathbb{R}^d} dw \sqrt{2-2\cos\theta_{d-2}} \delta(w-1) \\ &= c_2^d \int_0^\pi d\varphi \int_0^\pi d\theta_1 \dots \int_0^\pi d\theta_{d-2} v_2 \sqrt{1-\cos\theta_{d-2}} \prod_{k=1}^{d-2} (\sin\theta_k)^k \\ &= v_2 c_2^d 2\pi \left(\prod_{k=1}^{d-2} \int_0^\pi d\theta_k (\sin\theta_k)^k \right) \int_0^\pi d\theta_{d-2} (\sin\theta_{d-2})^{d-2} \sqrt{1-\cos\theta_{d-2}} \\ &= c_2^d v_2 2\pi J_d F_d = J_d = F_d \end{aligned}$$

Alors:

$$\begin{cases} \dot{X} = -\sigma^{d-1} \beta_d \left[X^2 \cdot 0 + \frac{XY}{\int_{\mathbb{R}^d} \delta(v_2-c_2)} c_2^d \int_{\mathbb{R}^d} \delta(v_2-c_2) \right] = -\sigma^{d-1} \beta_d c_2 XY \\ \dot{Y} = -\sigma^{d-1} \beta_d \left[\frac{Y^2}{\int_{\mathbb{R}^d} \delta(v_2-c_2)} c_2^d v_2 2\pi J_d F_d + XY c_2 \right] = -\sigma^{d-1} \beta_d \left[\frac{Y^2 c_2 v_2 2\pi J_d F_d}{\int_{\mathbb{R}^d} \delta(v_2-c_2)} + XY c_2 \right] \end{cases}$$

$$\Rightarrow \begin{cases} \dot{X} = -\sigma^{d-1} \beta_d c_2 XY \\ \dot{Y} = -\sigma^{d-1} \beta_d c_2 \left[Y^2 v_2 \frac{2\pi J_d F_d}{\int_{\mathbb{R}^d} \delta(v_2-c_2)} + XY \right] \end{cases}$$

$$\Rightarrow \begin{cases} \dot{X} = -\sigma^{d-1} \beta_d c_2 \frac{v_2 2\pi J_d F_d}{\int_{\mathbb{R}^d} \delta(v_2-c_2)} \cdot \frac{\int_{\mathbb{R}^d} \delta(v_2-c_2)}{v_2 2\pi J_d F_d} XY \\ \dot{Y} = -\sigma^{d-1} \beta_d c_2 \frac{v_2 2\pi J_d F_d}{\int_{\mathbb{R}^d} \delta(v_2-c_2)} \left[Y^2 + XY \frac{\int_{\mathbb{R}^d} \delta(v_2-c_2)}{v_2 2\pi J_d F_d} \right] \end{cases}$$

$$\Rightarrow \begin{cases} \dot{X} = -XY \frac{\int_{\mathbb{R}^d} \delta(v_2-c_2)}{v_2 2\pi J_d F_d} \\ \dot{Y} = -Y^2 - XY \frac{\int_{\mathbb{R}^d} \delta(v_2-c_2)}{v_2 2\pi J_d F_d} \end{cases}$$

$\tau = \epsilon \sigma^{d-1} 2\pi \frac{\beta_d \int_{\mathbb{R}^d} \delta(v_2-c_2)}{\int_{\mathbb{R}^d} \delta(v_2-c_2)} v_2 c_2$
 même temps que dans le cas général
 $C_1 \neq 0$: OK.

Soit $F_0 := \frac{J_d}{\sqrt{2\pi J_d F_0}}$, alors :

$$\begin{cases} \dot{X} = -XY F_0 \\ \dot{Y} = -Y^2 - XY F_0 \end{cases}$$

Pour achever la preuve, comme les temps τ sont les mêmes et ne dépendent pas de C_1 ni de C_2 , alors il suffit de vérifier que

$\lim_{q \rightarrow 0} F = F_0$

i.e. $\lim_{C_1 \rightarrow 0} \frac{\sqrt{C_1^2 + C_2^2}}{\sqrt{2} C_2} \frac{\int_{-1}^1 dx \sqrt{1 - \frac{2 C_1 C_2}{C_1^2 + C_2^2} x} / (1-x^2)^{\frac{d-3}{2}}}{\int_{-1}^1 dx \sqrt{1-x} / (1-x^2)^{\frac{d-3}{2}}} = \frac{2\pi^{d/2}}{\Gamma(d/2)} \frac{1}{\sqrt{2} 2\pi} \frac{1}{F_d} \frac{1}{J_d}$

$\Rightarrow \frac{1}{\sqrt{2}} \int_{-1}^1 dx \frac{1}{(1-x^2)^{\frac{d-3}{2}}} = \frac{\cancel{2\pi}^{d/2}}{\Gamma(d/2)} \frac{1}{\cancel{2\pi}} \frac{1}{\prod_{k=1}^{\frac{d-3}{2}} \left(\frac{\Gamma(k+1)}{\Gamma(\frac{k+1}{2})} \right) \Gamma(1+k/2)}$; $\pi^{\frac{d}{2} - \frac{3}{2} - \frac{d-3}{2}} = \pi^{\frac{1}{2}(d-2-\frac{d-3}{2})} = \pi^{1/2}$

$\Rightarrow \int_{-1}^1 dx \frac{1}{(1-x^2)^{\frac{d-3}{2}}} = \frac{\sqrt{\pi}}{\Gamma(\frac{d}{2}) \prod_{k=1}^{\frac{d-3}{2}} \Gamma(\frac{k+1}{2}) \Gamma(1+k/2)}$

ce qui est a priori bien d'être évident. On admet que le cas général est correct.

Comportement asymptotique :

$$\frac{dY}{dX} = \frac{Y^2 + XY F}{\delta X^2 + XY F} = \frac{Y^2/X^2 + Y/X F}{\delta + Y/X F} = \frac{V^2 + VF}{\delta + VF}$$

$\Rightarrow V + X \frac{dV}{dX} = \frac{V^2 + VF}{\delta + VF}$

$\Rightarrow X \frac{dV}{dX} = \frac{V^2 + VF - \delta V - V^2 F}{\delta + VF} = \frac{V^2(1-F) + V(F-\delta)}{\delta + VF} = \frac{[V(1-F) + F - \delta] V}{\delta + VF}$

$\Rightarrow \frac{dX}{X} = dV \frac{\delta + VF}{[V(1-F) + F - \delta] V} = dV \left[\frac{\alpha}{V} + \frac{\beta'}{V(1-F) + F - \delta} \right] = dV \left[\frac{\alpha [V(1-F) + F - \delta] + \beta' V}{V [V(1-F) + F - \delta]} \right]$

$\Rightarrow \alpha [V(1-F) + F - \delta] + \beta' V = \delta + VF$

$\Rightarrow \begin{cases} \alpha(1-F) + \beta' = F \\ \alpha(F-\delta) = \delta \end{cases} \Rightarrow \beta' = F - \frac{\delta(1-F)}{F-\delta}$

$\Rightarrow \frac{dX}{X} = dV \left[\frac{\delta}{F-\delta} \frac{1}{V} + \underbrace{\left(F - \frac{\delta(1-F)}{F-\delta} \right)}_{:= \beta} \frac{1}{1-F} \frac{1}{V + \frac{F-\delta}{1-F}} \right]$

$\Rightarrow \frac{dX}{X} = dV \left[\frac{\alpha}{V} + \frac{\beta}{V + \frac{F-\delta}{1-F}} \right]$; $\alpha = \frac{\delta}{F-\delta} > 0$; $\beta = \frac{F}{1-F} - \frac{\delta}{F-\delta} = \frac{F}{1-F} - \alpha > 0$

$\Rightarrow \ln \left(\frac{X}{X_0} \right) = \alpha \ln \left(\frac{V}{V_0} \right) + \beta \ln \left(\frac{V+\mu}{V_0+\mu} \right)$; $\mu = \frac{F-\delta}{1-F}$

$\Rightarrow \left(\frac{X}{X_0} \right) = \left(\frac{V}{V_0} \right)^\alpha \left(\frac{V+\mu}{V_0+\mu} \right)^\beta$; $\mu = \frac{F-\delta}{1-F}$; $\alpha = \frac{\delta}{F-\delta} > 0$; $\beta = \frac{F}{1-F} - \alpha > 0$

Car limiter :

$\lim_{\tau \rightarrow \infty} \frac{X}{X_0} = \frac{X_\infty}{X_0} = 0$
 $\lim_{C_1 \rightarrow 0} \frac{X}{X_0} = \left(\frac{V+\mu_0}{V_0+\mu_0} \right)^{\beta/(1-F_0)}$; $\mu_0 = \frac{F_0}{1-F_0}$
 $= \left(\frac{V+1/k}{V_0+1/k} \right)^{1/k}$; $k = \frac{1-F_0}{F_0}$

: m. structure qu'on \mathbb{Z}^d : ok.
 ↳ mais avec redéf. de \mathcal{H} .

Equation implicite pour $V(\tau)$:

$$\frac{d}{d\tau} \left(\frac{1}{X} \right) = -\frac{1}{X^2} \dot{X} = -\frac{1}{X^2} [-X^2 \delta - X Y F] = \delta + \frac{Y}{X} F$$

$$\Rightarrow \frac{d}{d\tau} \left(\frac{X_0}{X} \right) = \delta X_0 + X_0 V F$$

$$\Rightarrow \frac{dV}{d\tau} \frac{d}{dV} \left(\frac{X_0}{X} \right) = \delta X_0 + X_0 V F$$

$$\Rightarrow \frac{dV}{d\tau} \frac{d}{dV} \left(\frac{X_0}{X} \right) = (\delta X_0 + X_0 V F) d\tau$$

$$\Rightarrow X_0 d\tau = \frac{1}{\delta + V F} \frac{d}{dV} \left(\frac{X_0}{X} \right)$$

$$\Rightarrow X_0 \tau = \int_V^{V_0} \frac{1}{\delta + F \cdot s} \frac{d}{ds} \left[- \left(\frac{V_0}{s} \right)^\alpha \left(\frac{V_0 + \mu}{s + \mu} \right)^\beta \right]$$

; $\alpha > 0; \beta > 0$

Comportement asymptotique pour $X(\tau)$:

$$\frac{d}{d\tau} \left(\frac{1}{X} \right) = \delta + V F \quad ; \quad \lim_{\tau \rightarrow \infty} V(\tau) = 0$$

$$\Rightarrow d \left(\frac{1}{X} \right) \approx \delta d\tau$$

$$\Rightarrow \frac{1}{X} - \frac{1}{X_0} \approx \delta \tau$$

$$\Rightarrow \frac{X(\tau)}{X_0} \approx \frac{1}{1 + \delta X_0 \tau}$$

Borne: a) $\frac{1}{\delta} \gg \frac{1}{\delta + F \cdot s} \Rightarrow X_0 \tau \leq \frac{1}{\delta} \left[\frac{X_0}{X} - 1 \right] \Rightarrow X(\tau) \leq \frac{X_0}{1 + \delta X_0 \tau}$

b) $\frac{1}{\delta + F \cdot s} \gg \frac{1}{\delta + F V_0} \Rightarrow X_0 \tau \geq \frac{1}{\delta + F V_0} \left[\frac{X_0}{X} - 1 \right] \Rightarrow X(\tau) \geq \frac{X_0}{1 + (\delta + F V_0) \tau}$

$$\Rightarrow \frac{X_0}{1 + (\delta + F V_0) \tau} \leq X(\tau) \leq \frac{X_0}{1 + \delta X_0 \tau}$$

Comportement asymptotique pour $Y(\tau)$:

$$\frac{d}{d\tau} \left(\frac{1}{Y} \right) = -\frac{1}{Y^2} \dot{Y} = -\frac{1}{Y^2} (-Y^2 - X Y F) = 1 + \frac{X}{Y} F = 1 + \frac{1}{V} F \quad (*)$$

Comp. asymp. pour $V(\tau)$:

$$X_0 \tau = \int_V^{V_0} \frac{1}{\delta + F \cdot s} \frac{d}{ds} \left[- \left(\frac{V_0}{s} \right)^\alpha \left(\frac{V_0 + \mu}{s + \mu} \right)^\beta \right]$$

$$\approx \frac{1}{\delta} \left(\frac{V_0}{\mu} + 1 \right)^\beta \left[- \left(\frac{V_0}{s} \right)^\alpha \right]_V^{V_0}$$

$$\approx \frac{1}{\delta} \left(\frac{V_0}{\mu} + 1 \right)^\beta \left(\frac{V_0}{V} \right)^\alpha$$

$$\Rightarrow V^\alpha \approx \frac{1}{\delta} \left(\frac{V_0}{\mu} + 1 \right)^\beta \frac{1}{X_0 \tau}$$

$$\Rightarrow V(\tau) \approx \frac{1}{\delta^{1/\alpha}} \left(\frac{V_0}{\mu} + 1 \right)^{\beta/\alpha} (X_0 \tau)^{-1/\alpha}$$

demi (*) \Rightarrow

$$\frac{d}{d\tau} \left(\frac{1}{Y} \right) \approx \delta^{1/\alpha} \left(\frac{V_0}{\mu} + 1 \right)^{-\beta/\alpha} (X_0 \tau)^{1/\alpha} F$$

$$\Rightarrow \frac{1}{Y} - \frac{1}{Y_0} \approx \delta^{1/\alpha} \left(\frac{V_0}{\mu} + 1 \right)^{-\beta/\alpha} X_0^{1/\alpha} \frac{1}{1/\alpha + 1} \tau^{1/\alpha + 1} F$$

$$\Rightarrow Y(\tau) \approx \delta^{-1/\alpha} \left(\frac{v_0}{\lambda} + 1\right)^{\beta/\alpha} x_0^{-1/\alpha} \left(\frac{1}{\alpha} + 1\right) \tau^{-(1/\alpha + 1)} F^{-1} \quad (6)$$

avec:

$$\frac{1}{\alpha} + 1 = \frac{F - \delta}{\delta} + \frac{\delta}{\delta} = \frac{F}{\delta}$$

$$\begin{aligned} \Rightarrow Y(\tau) &\approx \delta^{-1/\alpha} \left(\frac{v_0}{\lambda} + 1\right)^{\beta/\alpha} x_0^{-1/\alpha} \frac{F}{\delta} \frac{1}{F} \tau^{-F/\delta} \\ &\approx \delta^{-(1/\alpha + 1)} \left(\frac{v_0}{\lambda} + 1\right)^{\beta/\alpha} \frac{1}{x_0^{1/\alpha}} \tau^{-F/\delta} \end{aligned}$$

$$\Rightarrow \boxed{Y(\tau) \approx \frac{1}{\delta^{F/\delta} x_0^{1/\alpha}} \left(\frac{v_0}{\lambda} + 1\right)^{\beta/\alpha} \tau^{-F/\delta}}$$

$$\mu = \frac{F - \delta}{1 - F} ; \delta = c/2 ; \alpha = \frac{\delta}{F - \delta} > 0 ; \beta = \frac{F}{1 - F} - \alpha > 0$$

Même structure que $d=2, 3, \dots$

Dimension arbitraire $d > 3$: cas particulier $c_1 = 0$

On a déjà établi: (pages 3,4)

$$\begin{cases} \dot{X} = -XYF_0 \\ \dot{Y} = -Y^2 - XYF_0 \end{cases}$$

An nouveau, mêmes calculs que pour $d=2,3$.

Résolution:

$$\frac{dY}{dX} = \frac{Y^2 + XYF_0}{XYF_0} = 1 + \frac{Y}{X} \frac{1}{F_0} = 1 + \frac{1}{F_0} V$$

$$\Rightarrow V + X \frac{dV}{dX} = 1 + \frac{1}{F_0} V$$

$$\Rightarrow X \frac{dV}{dX} = 1 + \underbrace{\left(\frac{1}{F_0} - 1\right)}_{=: k} V \quad ; \text{ même éq. que } d=2, \text{ avec différent } k.$$

$$\Rightarrow \frac{dX}{X} = \frac{dV}{1+kV}$$

$$\Rightarrow \ln\left(\frac{X}{X_0}\right) = \frac{1}{k} \ln\left(\frac{1+kV}{1+kV_0}\right)$$

$$\Rightarrow \frac{X}{X_0} = \left(\frac{1+kV}{1+kV_0}\right)^{1/k}$$

Calcul de X_0 :

$$\lim_{\tau \rightarrow \infty} \left(\frac{X}{X_0}\right) = \frac{X_\infty}{X_0} = \frac{1}{(1+kV_\infty)^{1/k}}$$

$$\Rightarrow X_\infty = \frac{X_0}{(1+kV_\infty)^{1/k}}$$

Comportement asymptotique

$$\frac{d}{d\tau} \left(\frac{1}{X}\right) = -\frac{1}{X^2} \dot{X} = -\frac{1}{X^2} (-XYF_0) = \frac{Y}{X} F_0 = VF_0$$

$$\Rightarrow \frac{d}{d\tau} \left(\frac{X_0}{X}\right) = V X_0 F_0 \quad ; \frac{d}{d\tau} = \frac{dV}{dV} \frac{dV}{d\tau}$$

$$\Rightarrow \frac{dV}{d\tau} \frac{d}{dV} \left(\frac{X_0}{X}\right) = V X_0 F_0$$

$$\Rightarrow F_0 X_0 d\tau = dV \frac{1}{V} \frac{d}{dV} \left[\left(\frac{1+kV}{1+kV_0}\right)^{1/k} \right]$$

$$\Rightarrow F_0 X_0 \tau = \int_{V_0}^V \frac{1}{s} \frac{d}{ds} \left[-\left(\frac{1+kV_0}{1+k+s}\right)^{1/k} \right]$$

Intégration par parties: c'est exactement la même expression que le cas 2d, $c_1 = 0$, eq. (10), avec $x_0 := x_0 F_0$, donc

$$F_0 X_0 \tau \simeq -I - \ln\left(\frac{V(1+kV_0)^{1/k}}{V_0^{1/(1+kV_0)}}\right) \quad ; \quad I = \int_0^{V_0} ds h(s) \frac{d^2}{ds^2} \left[-\left(\frac{1+kV_0}{1+k+s}\right)^{1/k} \right]$$

$$\Rightarrow V(\tau) \simeq V_0 \frac{1}{(1+kV_0)^{1/k+1}} \exp\left(-\frac{I}{(1+kV_0)^{1/k}}\right) \exp\left(-F_0 \frac{X_0}{(1+kV_0)^{1/k}} \tau\right)$$

$= I \frac{X_\infty}{X_0} \quad \quad \quad = X_\infty$

$$\Rightarrow V(\tau) \simeq f e^{-F_0 X_\infty \tau} \quad ; \quad f = V_0^{1/(1+kV_0)^{1/k+1}} \cdot e^{-I \frac{X_\infty}{X_0}}$$

Comportements asymptotiques de X et Y

$$\frac{d}{d\tau} \left(\frac{1}{X}\right) = VF_0 = f F_0 e^{-F_0 X_\infty \tau} \quad ; \quad \int_\tau^\infty$$

$$\Rightarrow \frac{1}{X_\infty} - \frac{1}{X(\tau)} = f F_0 \frac{1}{-F_0 X_\infty} e^{-F_0 X_\infty s} \Big|_\tau^\infty = \frac{f}{X_\infty} e^{-F_0 X_\infty \tau}$$

$$\Rightarrow \frac{1}{X(\tau)} = \frac{1 - f e^{-F_0 X_\infty \tau}}{X_\infty}$$

$$\Rightarrow X(\tau) = X_\infty \frac{1}{1 - f e^{-F_0 X_\infty \tau}} \approx X_\infty [1 + f e^{-F_0 X_\infty \tau}]$$

pour $Y(\tau)$:

$$\dot{X} = -XY F_0$$

$$\Rightarrow \cancel{X_\infty} [0 + f(+f_0 X_\infty) e^{-F_0 X_\infty \tau}] = + \cancel{X_\infty} [1 + f e^{-F_0 X_\infty \tau}] Y \cancel{f_0}$$

$0(e^{-2F_0 X_\infty \tau})$

$$\Rightarrow Y(\tau) \approx f X_\infty e^{-F_0 X_\infty \tau}$$

Relation entre X et Y :

$$X(\tau) = X_\infty + Y(\tau)$$

Equation:

$$\frac{\partial}{\partial t} f(v_1, t) = -2\sigma f(v_1, t) \int_{\mathbb{R}^2} dv_2 |v_2| f(v_2, t) \quad (1)$$

Condition initiale (distribution):

$$f(v, t) = X \frac{1}{2\pi c_1} \delta(v - c_1) + Y \frac{1}{2\pi c_2} \delta(v - c_2) \quad ; \quad c_1 \neq c_2 \quad ; \quad c_1 > 0 \quad ; \quad c_2 > c_1 \quad (2)$$

Eqn. pour X et Y: (2) dans (1) \Rightarrow

$$\dot{X} \frac{1}{2\pi c_1} \delta(v - c_1) + \dot{Y} \frac{1}{2\pi c_2} \delta(v - c_2) = -2\sigma \left[X \frac{1}{2\pi c_1} \delta(v - c_1) + Y \frac{1}{2\pi c_2} \delta(v - c_2) \right] \int_{\mathbb{R}^2} dv_2 |v_2| \left[\frac{X}{2\pi c_1} \delta(v_2 - c_1) + \frac{Y}{2\pi c_2} \delta(v_2 - c_2) \right] \quad (3)$$

$$\stackrel{(3)}{\Rightarrow} \dot{X} \frac{1}{2\pi c_1} = -2\sigma X \frac{1}{2\pi c_1} \left[\underbrace{\frac{X}{2\pi c_1} \int_{\mathbb{R}^2} dv_2 |v_2| \delta(v_2 - c_1)}_{:= I_{11}} + \underbrace{\frac{Y}{2\pi c_2} \int_{\mathbb{R}^2} dv_2 |v_2| \delta(v_2 - c_2)}_{:= I_{21}} \right] \quad ; \quad |v_1| = c_1 \quad (4)$$

$$\stackrel{(3)}{\Rightarrow} \dot{Y} \frac{1}{2\pi c_2} = -2\sigma Y \frac{1}{2\pi c_2} \left[\underbrace{\frac{X}{2\pi c_1} \int_{\mathbb{R}^2} dv_2 |v_2| \delta(v_2 - c_1)}_{:= I_{12}} + \underbrace{\frac{Y}{2\pi c_2} \int_{\mathbb{R}^2} dv_2 |v_2| \delta(v_2 - c_2)}_{:= I_{22}} \right] \quad ; \quad |v_1| = c_2 \quad (5)$$

Calcul des I_{ij} : interprétation: I_{ij} = intégrale avec $|v_1| = c_j$ et $|v_2| = c_i$

$$\begin{aligned} I_{ij} &= \int_{\mathbb{R}^2} dv_2 |v_2| \delta(v_2 - c_i) \Big|_{|v_1|=c_j} = \int_{\mathbb{R}^2} dv_2 |\hat{v}_1 c_j - v_2| \delta(v_2 - c_i) \quad ; \quad w = v_2 / c_i \\ &= \int_{\mathbb{R}^2} dw c_i^2 |\hat{v}_1 c_j - \hat{w} c_i| \delta(w - 1) \frac{1}{c_i} \\ &= c_i \int_{\mathbb{R}^2} dw |\hat{v}_1 c_j - \hat{w} c_i| \delta(w - 1) \end{aligned} \quad (6)$$

$$\begin{aligned} 1) \underline{i=j}: I_{ii} &\stackrel{(6)}{=} c_i^2 \int dw |\hat{v}_1 - \hat{w}| \delta(w - 1) \\ &= c_i^2 \int_0^{2\pi} d\theta \sqrt{\hat{v}_1^2 + \hat{w}^2 - 2\hat{v}_1 \hat{w} \cos\theta} \quad ; \quad \hat{v}_1^2 = \hat{w}^2 = 1 \\ &= c_i^2 \int_0^{2\pi} d\theta \sqrt{2 - 2\cos\theta} \quad ; \quad 1 - \cos\theta = 2(\sin\frac{\theta}{2})^2 \\ &= 2c_i^2 \int_0^{2\pi} d\theta |\sin\frac{\theta}{2}| \quad ; \quad x = \theta/2 \quad ; \quad dx = 1/2 d\theta \\ &= 4c_i^2 \int_0^\pi dx \sin(x) \\ &= 8c_i^2 \end{aligned} \quad (7)$$

$$\begin{aligned} 2) \underline{i \neq j}: I_{ij} &\stackrel{(6)}{=} c_i \int_0^{2\pi} d\theta \sqrt{c_i^2 + c_j^2 - 2c_i c_j \cos\theta} \\ &= -2c_i c_j + 2c_i c_j - 2c_i c_j \cos\theta = -2c_i c_j + 2c_i c_j \underbrace{(1 - \cos\theta)}_{= 2(\sin\frac{\theta}{2})^2} \\ &= c_i \int_0^{2\pi} d\theta \sqrt{c_i^2 + c_j^2 - 2c_i c_j + 4c_i c_j (\sin\frac{\theta}{2})^2} \quad ; \quad x = \theta/2 \quad ; \quad dx = d\theta/2 \\ &= 2c_i |c_i - c_j| \int_0^\pi dx \sqrt{1 + \frac{4c_i c_j}{(c_i - c_j)^2} \sin^2 x} \\ &= 4c_i |c_i - c_j| E[k] \quad ; \quad E[k] = \int_0^{\pi/2} dx \sqrt{1 + k^2 \sin^2 x} \\ & \quad \quad \quad k = 2\sqrt{c_i c_j} / |c_i - c_j| \\ & \quad \quad \quad \lim_{c_i \rightarrow c_j} E[k] = \lim_{c_j \rightarrow 0} E[k] = \frac{\pi}{2} \end{aligned} \quad (8)$$

Eqn. pour X et Y: (4) et (5) \Rightarrow

$$\boxed{\begin{aligned} \dot{X} &= -\frac{\sigma}{\pi} \left[X^2 \frac{I_{11}}{c_1} + XY \frac{I_{21}}{c_2} \right] \\ \dot{Y} &= -\frac{\sigma}{\pi} \left[Y^2 \frac{I_{22}}{c_2} + XY \frac{I_{12}}{c_1} \right] \end{aligned}} \quad (9)$$

Expressions explicites: (7), (8) et (9) \Rightarrow

$$\begin{cases} \dot{X} = -\frac{\sigma}{\pi} \left[X^2 \frac{8c_1^2}{c_1} + XY \frac{4c_2 |c_1 - c_2| E[k]}{c_2} \right] & \left(\frac{2}{\pi} \right) \left(\frac{\pi}{2} \right) \cdot \left(\frac{c_2}{c_1} \right) \\ \dot{Y} = -\frac{\sigma}{\pi} \left[Y^2 \frac{8c_2^2}{c_2} + XY \frac{4c_1 |c_1 - c_2| E[k]}{c_1} \right] & \left(\frac{2}{\pi} \right) \left(\frac{\pi}{2} \right) \cdot \left(\frac{c_1}{c_2} \right) \end{cases}$$

$$\Rightarrow \begin{cases} \dot{X} = -\frac{\sigma}{\pi} \cdot \frac{\pi}{2} \cdot c_2 \left[X^2 \frac{8c_1}{c_2} \cdot \frac{2}{\pi} + XY \frac{4|c_1-c_2|}{c_2} E[k] \frac{2}{\pi} \right] \\ \dot{Y} = -\frac{\sigma}{\pi} \cdot \frac{\pi}{2} \cdot c_2 \left[Y^2 \frac{8c_2}{c_2} \cdot \frac{2}{\pi} + XY \frac{4|c_1-c_2|}{c_2} E[k] \frac{2}{\pi} \right] \end{cases}$$

$$\delta = \frac{c_1}{c_2}$$

$$\Rightarrow \begin{cases} \dot{X} = -2\sigma c_2 \left[X^2 \delta \frac{4}{\pi} + XY \frac{|c_1-c_2|}{c_2} E[k] \frac{2}{\pi} \right] \\ \dot{Y} = -2\sigma c_2 \left[Y^2 \frac{4}{\pi} + XY \frac{|c_1-c_2|}{c_2} E[k] \frac{2}{\pi} \right] \end{cases}$$

$$\tau = 2\sigma c_2 t$$

$$F(c_1, c_2) = \frac{|c_1-c_2|}{c_2} E[k] \frac{2}{\pi}$$

$$\Rightarrow \begin{cases} \dot{X} = -\delta \frac{4}{\pi} X^2 - XY F(c_1, c_2) \\ \dot{Y} = -\frac{4}{\pi} Y^2 - XY F(c_1, c_2) \end{cases}$$

$$F(c_1, c_2) = \frac{2}{\pi} \frac{|c_1-c_2|}{c_2} \int_0^{\pi/2} dx \sqrt{1 + \frac{4c_1 c_2}{|c_1-c_2|^2} \sin^2 x} = \frac{2}{\pi} \int_0^{\pi/2} dx \sqrt{\frac{c_1^2 \cos^2 x + c_2^2 \sin^2 x}{c_2^2}} \quad (10)$$

$$= \frac{2}{\pi} \int_0^{\pi/2} dx \sqrt{\frac{c_1^2}{c_2^2} \cos^2 x + \frac{4c_1}{c_2} \sin^2 x}$$

avec: $F(c_1, c_2) = \frac{|c_1-c_2|}{c_2} E[k] \frac{2}{\pi}$; $\lim_{c_1 \rightarrow 0} F(c_1, c_2) = 1$

$$E[k] = \int_0^{\pi/2} dx \sqrt{1+k^2 \sin^2 x}$$
 ; $k = 2 \sqrt{c_1 c_2} / |c_1 - c_2|$; $\delta = c_1/c_2$

Dans la limite $c_1 \rightarrow 0$ on retrouve bien les équations établies dans le cas particulier $c_1 = 0$.

Comportement asymptotique: (10) \Rightarrow

$$\frac{dY}{dX} = \frac{Y^2 \frac{4}{\pi} + XY F}{\delta X^2 \frac{4}{\pi} + XY F} = \frac{Y^2/X^2 \frac{4}{\pi} + X/Y F}{\delta \frac{4}{\pi} + X/Y F} = \frac{V^2 \frac{4}{\pi} + VF}{\delta \frac{4}{\pi} + VF}$$

$$\Rightarrow V + X \frac{dV}{dX} = \frac{V^2 \frac{4}{\pi} + VF}{\delta \frac{4}{\pi} + VF}$$

$$\Rightarrow X \frac{dV}{dX} = \frac{V^2 \frac{4}{\pi} + VF - V \frac{4}{\pi} \delta - V^2 F}{\delta \frac{4}{\pi} + VF} = \frac{V^2 (\frac{4}{\pi} - F) + V(F - \frac{4}{\pi} \delta)}{\delta \frac{4}{\pi} + VF}$$

$$\Rightarrow \frac{dX}{X} = dV \frac{\delta \frac{4}{\pi} + VF}{V^2 (\frac{4}{\pi} - F) - V(\frac{4}{\pi} \delta - F)} = dV \frac{\delta \frac{4}{\pi} + VF}{[V(\frac{4}{\pi} - F) - \frac{4}{\pi} \delta + F] V}$$

$$= dV \left[\frac{\alpha}{V} + \frac{\beta'}{V(\frac{4}{\pi} - F) - \frac{4}{\pi} \delta + F} \right] = dV \left[\frac{\alpha [V(\frac{4}{\pi} - F) - \frac{4}{\pi} \delta + F] + \beta'}{[V(\frac{4}{\pi} - F) - \frac{4}{\pi} \delta + F] V} \right] \quad (11)$$

$$\Rightarrow \alpha V(\frac{4}{\pi} - F) - \alpha \frac{4}{\pi} \delta + \alpha F + \beta' = \delta \frac{4}{\pi} + VF$$

$$\Rightarrow \begin{cases} \alpha (\frac{4}{\pi} - F) + \beta' = F & \Rightarrow \beta' = F - \alpha (\frac{4}{\pi} - F) \\ -\alpha \frac{4}{\pi} \delta + \alpha F = \delta \frac{4}{\pi} & \Rightarrow \alpha (F - \frac{4}{\pi} \delta) = \delta \frac{4}{\pi} \Rightarrow \alpha = \delta \frac{4}{\pi} / (F - \frac{4}{\pi} \delta) \end{cases} \quad (12)$$

$$\Rightarrow \beta' = F - \frac{\delta \frac{4}{\pi}}{F - \frac{4}{\pi} \delta} (\frac{4}{\pi} - F)$$

(12) dans (11) \Rightarrow

$$\frac{dX}{X} = dV \left(\frac{\alpha}{V} + \frac{\beta'}{V(\frac{4}{\pi} - F) - \frac{4}{\pi} \delta + F} \right) \quad ; \quad \beta = \frac{F}{\frac{4}{\pi} - F} - \frac{\delta \frac{4}{\pi}}{F - \frac{4}{\pi} \delta}$$

$$\Rightarrow \ln\left(\frac{X}{X_0}\right) = \alpha \ln\left(\frac{V}{V_0}\right) + \beta \ln\left(\frac{V + \frac{F - \delta \frac{4}{\pi}}{\frac{4}{\pi} - F}}{V_0 + \frac{F - \delta \frac{4}{\pi}}{\frac{4}{\pi} - F}}\right)$$

$$\Rightarrow \frac{X}{X_0} = \left(\frac{V}{V_0}\right)^\alpha \left(\frac{V + \frac{F - \delta \frac{4}{\pi}}{\frac{4}{\pi} - F}}{V_0 + \frac{F - \delta \frac{4}{\pi}}{\frac{4}{\pi} - F}}\right)^\beta \quad (13)$$

avec: $\alpha = \frac{\delta \frac{4}{\pi}}{F - \delta \frac{4}{\pi}} \rightarrow 0$; $\beta = \frac{F}{\frac{4}{\pi} - F} - \alpha > 0$; $\lim_{c_1 \rightarrow 0} \alpha = 0$; $\lim_{c_1 \rightarrow 0} \beta = \frac{1}{\frac{4}{\pi} - 1} = \frac{1}{e}$

Cas limite:

$$\lim_{c_1 \rightarrow 0} \frac{X}{X_0} = \frac{X_{\infty}}{X_0} = 0 \quad ; \quad \text{car: par d'espèce survivante dès que } c_1 > 0$$

$$\lim_{c_1 \rightarrow 0} \frac{X}{X_0} = \left(\frac{V + 1/(4\pi - 2)}{V_0 + 1/(4\pi - 2)}\right)^{1/e} = \left(\frac{V + 1/e}{V_0 + 1/e}\right)^{1/e} \quad ; \quad \text{or, le m. résultat déjà obtenu pour } c_1 = 0.$$

Equation implicite pour $V(\tau)$:

$$\frac{d}{d\tau} \left(\frac{1}{X} \right) = -\frac{1}{X^2} \dot{X} \stackrel{(10)}{=} -\frac{1}{X^2} (-\delta \frac{4}{\pi} X^2 - X Y F) = \delta \frac{4}{\pi} + \frac{Y}{X} F = \delta \frac{4}{\pi} + V F \quad \Rightarrow \frac{d}{d\tau} = \frac{dV}{dV} \frac{dV}{d\tau}$$

$$\Rightarrow \frac{d}{d\tau} \left(\frac{X_0}{X} \right) = \left(\delta \frac{4}{\pi} + V F \right) X_0$$

$$\Rightarrow \frac{dV}{d\tau} \frac{d}{dV} \left(\frac{X_0}{X} \right) = \left(\delta \frac{4}{\pi} + V F \right) X_0$$

$$\Rightarrow X_0 d\tau \stackrel{(11)}{=} dV \frac{1}{\delta \frac{4}{\pi} + V F} \frac{d}{dV} \left(\left(\frac{V_0}{V} \right)^\alpha \left(\frac{V_0 + \frac{F - 4/\pi \delta}{4/\pi - F}}{V + \frac{F - 4/\pi \delta}{4/\pi - F}} \right)^\beta \right)$$

$$\Rightarrow X_0 \tau = \int_V^{V_0} ds \frac{1}{\delta \frac{4}{\pi} + F s} \frac{d}{ds} \left[- \left(\frac{V_0}{s} \right)^\alpha \left(\frac{V_0 + \frac{F - 4/\pi \delta}{4/\pi - F}}{s + \frac{F - 4/\pi \delta}{4/\pi - F}} \right)^\beta \right] \quad ; \alpha > 0; \beta > 0 \quad (14)$$

Comportement asymptotique pour $X(\tau)$: (10) \Rightarrow

$$\frac{d}{d\tau} \left(\frac{1}{X} \right) = \delta \frac{4}{\pi} + V F \quad , \quad \lim_{\tau \rightarrow \infty} V(\tau) = 0$$

$$\Rightarrow \frac{d}{d\tau} \left(\frac{1}{X} \right) \approx \frac{4}{\pi} \delta$$

$$\Rightarrow \frac{1}{X} - \frac{1}{X_0} \approx \frac{4}{\pi} \delta \tau$$

$$\Rightarrow X(\tau) \approx \frac{X_0}{1 + \frac{4\delta}{\pi} X_0 \tau} \quad (15)$$

Remarque: bornes: on peut aussi trouver des bornes sup. et inf. pour $X(\tau)$:

$$a) \frac{1}{\delta \frac{4}{\pi}} \geq \frac{1}{\delta \frac{4}{\pi} + F \cdot s} \Rightarrow X_0 \tau \leq \frac{\pi}{4\delta} \left[\frac{X_0}{X} - 1 \right] \Rightarrow X(\tau) \leq \frac{X_0}{1 + \frac{4\delta}{\pi} X_0 \tau} \quad (16)$$

$$b) \frac{1}{\delta \frac{4}{\pi} + F s} \geq \frac{1}{\delta \frac{4}{\pi} + F V_0} \Rightarrow X_0 \tau \geq \frac{1}{\delta \frac{4}{\pi} + F V_0} \left[\frac{X_0}{X} - 1 \right] \Rightarrow X(\tau) \geq \frac{X_0}{1 + \left(\frac{4\delta}{\pi} + F V_0 \right) \tau} \quad (17)$$

$$\Rightarrow \frac{X_0}{1 + \left(\frac{4\delta}{\pi} + F V_0 \right) \tau} \leq X(\tau) \leq \frac{X_0}{1 + \frac{4\delta}{\pi} X_0 \tau} \quad (18)$$

A nouveau, $X(\tau)$ converge vers la borne supérieure.

Comportement asymptotique pour $Y(\tau)$: (10) \Rightarrow

$$\frac{d}{d\tau} \left(\frac{1}{Y} \right) = -\frac{1}{Y^2} \dot{Y} \stackrel{(10)}{=} -\frac{1}{Y^2} \left[-\frac{4}{\pi} Y^2 - X Y F \right] = \frac{4}{\pi} + \frac{X}{Y} F = \frac{4}{\pi} + \sqrt{V} F \quad (19)$$

On doit d'abord déterminer le comportement asymptotique pour $V(\tau)$, avec (14):

$$X_0 \tau = \int_V^{V_0} ds \frac{1}{\delta \frac{4}{\pi} + F s} \frac{d}{ds} \left[- \left(\frac{V_0}{s} \right)^\alpha \left(\frac{V_0 + \mu}{s + \mu} \right)^\beta \right] \quad ; \mu = \frac{F - 4/\pi \delta}{4/\pi - F} \quad (20)$$

m. cas que βd : peut négliger s dans les termes ne produisant pas de divergence:

$$X_0 \tau \stackrel{\tau \rightarrow \infty}{\approx} \frac{\pi}{4\delta} \left(\frac{V_0 + \mu}{\mu} \right)^\beta \left[- \left(\frac{V_0}{s} \right)^\alpha \right]_V^{V_0} = -1 + \left(\frac{V_0}{V} \right)^\alpha \approx \left(\frac{V_0}{V} \right)^\alpha$$

$$\Rightarrow X_0 \tau V^\alpha \approx \frac{\pi}{4\delta} \left(\frac{V_0}{\mu} + 1 \right)^\beta V_0^\alpha$$

$$\Rightarrow V(\tau) \stackrel{\tau \rightarrow \infty}{\approx} V_0 \left(\frac{V_0}{\mu} + 1 \right)^{\beta/\alpha} \left(\frac{4\delta X_0}{\pi} \tau \right)^{-1/\alpha} \quad ; \mu = \frac{F - 4/\pi \delta}{4/\pi - F} \quad (21)$$

(21) dans (19) \Rightarrow

$$\frac{d}{d\tau} \left(\frac{1}{Y} \right) \stackrel{\tau \rightarrow \infty}{\approx} \frac{4}{\pi} + F \frac{1}{V_0} \left(\frac{V_0}{\mu} + 1 \right)^{-\beta/\alpha} \left(\frac{4\delta X_0}{\pi} \tau \right)^{1/\alpha}$$

négligeable devant $\tau^{1/\alpha}$

$$\frac{1}{Y} - \frac{1}{V_0} \xrightarrow[\tau \rightarrow \infty]{\text{négligeable}} F \frac{1}{V_0} \left(\frac{V_0}{M} + 1 \right)^{-\beta/\alpha} \left(\frac{4\delta X_0}{\pi} \right)^{1/\alpha} \frac{1}{\alpha + 1} \tau^{\frac{1}{\alpha} + 1}$$

$$\Rightarrow Y(\tau) \simeq \frac{V_0}{F} \left(\frac{V_0}{M} + 1 \right)^{\beta/\alpha} \left(\frac{4\delta}{\pi} \right)^{-1/\alpha} X_0^{-1/\alpha} \left(\frac{1}{\alpha} + 1 \right) \tau^{-(1/\alpha + 1)} \quad (22)$$

$$\text{avec: } \frac{1}{\alpha} + 1 = \frac{F - \delta^{4/\pi}}{\delta^{4/\pi}} + \frac{\delta^{4/\pi}}{\delta^{4/\pi}} = F \cdot \frac{\pi}{4\delta}$$

$$\Rightarrow Y(\tau) \simeq \frac{V_0}{F} \frac{\pi}{4\delta} \left(\frac{\pi}{4\delta} \right)^{1/\alpha} \left(\frac{V_0}{M} + 1 \right)^{\beta/\alpha} \tau^{-\frac{F\pi}{4\delta}} \frac{1}{X_0^{1/\alpha}}$$

$$= \left(\frac{\pi}{4\delta} \right)^{1/\alpha + 1}$$

$$= \left(\frac{\pi}{4\delta} \right)^{F^{4/\pi}}$$

$$\Rightarrow Y(\tau) \simeq \frac{V_0}{X_0^{1/\alpha}} \left(\frac{\pi}{4\delta} \right)^{\frac{\pi}{4\delta} F} \left(\frac{4/\pi - F}{F - 4/\pi\delta} V_0 + 1 \right)^{\beta/\alpha} \tau^{-\frac{\pi}{4\delta} F} \quad (23)$$

$$\alpha = \frac{\delta^{4/\pi}}{F - \delta^{4/\pi}} > 0; \quad \beta = \frac{F}{\frac{4}{\pi} - F} - \alpha > 0$$

Remarque: si $c_1 \rightarrow 0$; $c_1 > 0$, alors la décroissance de Y est d'autant plus rapide. En effet, dans ce cas on se rapproche du cas limite où $c_1 = 0$ pour lequel $Y(\tau)$ a une décroissance exponentielle, d'où la divergence de l'exposant de la loi de puissance qui "essaye d'approcher" l'exponentielle. On a bien entendu par là droit de réaliser formellement cette limite dans (23).

Car particulier $V_1 = 0$ dimension $d = 2$

Equation :

$$\begin{aligned} \frac{\partial}{\partial t} f(v_1, t) &= -\sigma^{d-1} \int_{\mathbb{R}^d} d\hat{\sigma} \theta(\hat{\sigma} \cdot v_{12}) (\hat{\sigma} \cdot \hat{v}_{12}) \int_{\mathbb{R}^d} dv_2 |v_{12}| f(v_1, t) f(v_2, t) \\ &\stackrel{d=2}{=} -\sigma \int_{\mathbb{R}^2} d\hat{\sigma} \theta(\hat{\sigma} \cdot v_{12}) (\hat{\sigma} \cdot \hat{v}_{12}) f(v_1, t) \int_{\mathbb{R}^2} dv_2 |v_{12}| f(v_2, t) \\ &= \pi^{1/2} \Gamma(3/2) \\ &= \pi^{1/2} / \Gamma(1+1/2) \\ &= \pi^{1/2} / \pi^{1/2} \\ &= 2 \end{aligned}$$

$$\Rightarrow \frac{\partial}{\partial t} f(v_1, t) = -2\sigma f(v_1, t) \int_{\mathbb{R}^2} dv_2 |v_{12}| f(v_2, t) \tag{1}$$

Condition initiale (distribution) :

$$f(v, t) = X(t) \delta(v) + Y(t) \frac{1}{2\pi c} \delta(v-c) \quad ; \quad c_1 = 0 ; c_2 = c \quad (\text{en effet : } \int_{\mathbb{R}^2} d\hat{\sigma} f(v, t) = x \int d\hat{\sigma} \delta(v) + \frac{1}{2\pi} \int d\hat{\sigma} \delta(v-c) = x + Y = n) \tag{2}$$

(2) dans (1) \Rightarrow

$$\dot{X} \delta(v) + \dot{Y} \frac{1}{2\pi c} \delta(v-c) = -2\sigma \left[X \delta(v) + Y \frac{1}{2\pi c} \delta(v-c) \right] \int_{\mathbb{R}^2} dv_2 |v_{12}| \left[X \delta(v_2) + Y \frac{1}{2\pi c} \delta(v_2-c) \right] \tag{3}$$

Coefficient $\delta(v_1)$:

$$\begin{aligned} X &= -2\sigma X \int_{\mathbb{R}^2} dv_2 |v_{12}| \left[X \delta(v_2) + \frac{Y}{2\pi c} \delta(v_2-c) \right] \\ &= -2\sigma X^2 \underbrace{\int_{\mathbb{R}^2} dv_2 |v_1 - v_2| \delta(v_2)}_{= 0 \text{ car } v_1 = v_2 = 0} - 2\sigma X Y \frac{1}{2\pi c} \int_{\mathbb{R}^2} dv_2 |v_1 - v_2| \delta(v_2 - c) \\ &= -2\sigma X Y c \end{aligned} \tag{4}$$

Coefficient $\delta(v_1 - c)$:

$$\begin{aligned} \dot{Y} \frac{1}{2\pi c} &= -2\sigma \frac{Y}{2\pi c} \int_{\mathbb{R}^2} dv_2 |v_{12}| \left[X \delta(v_2) + \frac{Y}{2\pi c} \delta(v_2 - c) \right] \\ &= -2\sigma \frac{XY}{2\pi c} \underbrace{\int_{\mathbb{R}^2} dv_2 |v_{12}| \delta(v_2)}_{= \int_{\mathbb{R}^2} dv_2 |v_1 - v_2| \delta(v_2)} - 2\sigma \frac{Y^2}{(2\pi c)^2} \int_{\mathbb{R}^2} dv_2 |v_{12}| \delta(v_2 - c) \\ &= \int_{\mathbb{R}^2} dv_2 |v_1 - v_2| \delta(v_2) = \int_{\mathbb{R}^2} dv_2 |c \hat{v}_1 - v_2| \delta(v_2) = c |\hat{v}_1| = c \\ &= \int_{\mathbb{R}^2} dw c^2 |c \hat{v}_1 - \hat{w}| \delta(cw - c) \quad ; \quad |v_2| = c \\ &= \int_{\mathbb{R}^2} dw c^2 |c \hat{v}_1 - \hat{w}| \delta(cw - c) \quad ; \quad w = v_2/c ; dw = \frac{1}{c^2} dv_2 \\ &= c^2 \int_0^{2\pi} d\theta \int_0^\infty dw |\hat{v}_1 - \hat{w}| \delta(w-1) \quad ; \quad |\hat{v}_1| = |\hat{w}| = 1 \\ &= c^2 \int_0^{2\pi} d\theta \sqrt{2-2\cos\theta} \quad ; \quad 1 - \cos\theta = 2(\sin(\theta/2))^2 \\ &= c^2 \sqrt{2} \int_0^{2\pi} d\theta \sqrt{2} \sin(\theta/2) \quad ; \quad x = \theta/2 ; dx = d\theta/2 \\ &= c^2 \cdot 2 \cdot \int_0^\pi dx 2 \sin x \\ &= 8c^2 \end{aligned}$$

$$= -2\sigma \frac{XY}{2\pi c} c - 2\sigma \frac{Y^2}{(2\pi c)^2} 8c^2$$

Conclusion :

$$\begin{cases} \dot{X} = -2\sigma c XY \\ \dot{Y} = -2\sigma c \left[XY + \frac{4}{\pi} Y^2 \right] \end{cases}$$

rescaling : $\tau = t \cdot 2\sigma c \Rightarrow$

$$\begin{cases} \dot{X} = -XY \\ \dot{Y} = -XY - \frac{4}{\pi} Y^2 \end{cases}$$

• Résolution: (5) ⇒

$$\frac{dY}{dX} = \frac{XY + 4/\pi Y^2}{XY} = 1 + \frac{4}{\pi} \frac{Y}{X} \quad ; \quad Y = VX \quad ; \quad Y' = V + XY'$$

$$\Rightarrow V + X \frac{dV}{dX} = 1 + \frac{4}{\pi} V$$

$$\Rightarrow X \frac{dV}{dX} = 1 + \frac{4(\pi-1)}{\pi} V$$

$$\Rightarrow \frac{dX}{X} = \frac{dV}{1 + \frac{4(\pi-1)}{\pi} V} = \frac{1}{\frac{4(\pi-1)}{\pi} V + 1} \frac{dV}{V}$$

$$\Rightarrow \ln\left(\frac{X}{X_0}\right) = \frac{1}{\frac{4(\pi-1)}{\pi}} \ln\left(\frac{\frac{4(\pi-1)}{\pi} V + 1}{\frac{4(\pi-1)}{\pi} V_0 + 1}\right)$$

$$\Rightarrow \boxed{\frac{X}{X_0} = \left(\frac{\frac{4(\pi-1)}{\pi} V + 1}{\frac{4(\pi-1)}{\pi} V_0 + 1}\right)^{\frac{\pi}{4(\pi-1)}}} \quad ; \quad X = \frac{4}{\pi} - 1 = 0.2732... \quad (6)$$

• Calcul de X_{∞} : (6) et $\lim_{\tau \rightarrow \infty} \Rightarrow$

$$\lim_{\tau \rightarrow \infty} \frac{X}{X_0} = \frac{X_{\infty}}{X_0} = \left(\frac{1/\pi}{1/\pi + V_0}\right)^{1/\pi} = \frac{1}{(1 + \pi V_0)^{1/\pi}} \quad (7)$$

$$\Rightarrow \boxed{X_{\infty} = \frac{X_0}{(1 + \pi V_0)^{1/\pi}}} \quad (8)$$

Si $X_0 = Y_0 = 1$, alors $X_{\infty} = (\pi/4)^{1/\pi} = 0.413095...$

• Comportement asymptotique:

$$(5) \Rightarrow \frac{d}{d\tau} \left(\frac{1}{X}\right) = -\frac{1}{X^2} \dot{X} \stackrel{(5)}{=} -\frac{1}{X^2} (-XY) = V \quad (9)$$

$$\Rightarrow \frac{d}{d\tau} \left(\frac{X_0}{X}\right) = X_0 V$$

$$\stackrel{(6)}{\Rightarrow} \frac{d}{d\tau} \left(\frac{1/\pi + V_0}{1/\pi + V}\right)^{1/\pi} = X_0 V \quad ; \quad \frac{d}{d\tau} = \frac{d}{dV} \frac{dV}{d\tau}$$

$$\Rightarrow \frac{1}{V} \frac{dV}{d\tau} \frac{d}{dV} \left(\frac{1/\pi + V_0}{1/\pi + V}\right)^{1/\pi} = X_0$$

$$\Rightarrow X_0 d\tau = \frac{1}{V} dV \frac{d}{dV} \left(\frac{1/\pi + V_0}{1/\pi + V}\right)^{1/\pi}$$

$$\Rightarrow X_0 \tau = \int_{V_0}^V ds \frac{1}{s} \frac{d}{ds} \left(\frac{1/\pi + V_0}{1/\pi + s}\right)^{1/\pi}$$

$$\Rightarrow \boxed{X_0 \tau = \int_V^{V_0} ds \frac{1}{s} \frac{d}{ds} \left[-\left(\frac{1/\pi + V_0}{1/\pi + s}\right)^{1/\pi}\right]} \quad (10)$$

Intégration par parties: $f = \ln(s) ; f' = 1/s$
 $g = \frac{d}{ds}[-\dots] ; g' = \frac{d^2}{ds^2}[-\dots]$

$$\Rightarrow X_0 \tau = \int_V^{V_0} ds f' g = -\int_V^{V_0} ds f g' + f g \Big|_V^{V_0}$$

$$= -\int_V^{V_0} ds \ln(s) \frac{d^2}{ds^2}[-\dots] + \ln(s) \frac{d}{ds} \left[-\left(\frac{1/\pi + V_0}{1/\pi + s}\right)^{1/\pi}\right] \Big|_V^{V_0}$$

approximation $\tau \rightarrow \infty$
 $\Rightarrow \int_V^{V_0} \rightarrow \int_0^{V_0}$

$$\approx -\int_0^{V_0} ds \ln(s) \frac{d^2}{ds^2} \left[-\left(\frac{1/\pi + V_0}{1/\pi + s}\right)^{1/\pi}\right] + \ln(s) \frac{1}{\pi} \frac{(1/\pi + V_0)^{1/\pi}}{(1/\pi + s)^{1/\pi+1}} \Big|_V^{V_0}$$

$$= \ln(V_0) \frac{1}{\pi} \frac{(1/\pi + V_0)^{1/\pi}}{(1/\pi + V_0)^{1/\pi+1}} - \ln(V) \frac{1}{\pi} \frac{(1/\pi + V_0)^{1/\pi}}{(1/\pi + V)^{1/\pi+1}}$$

$$= \frac{1}{\pi} \frac{1}{1/\pi + V_0} - \ln(V) \frac{1}{\pi} \frac{(1/\pi + V_0)^{1/\pi}}{(1/\pi + V)^{1/\pi+1}}$$

$$\approx -I + \ln(V_0) \frac{1}{1 + \pi V_0} - \ln(V) (1 + \pi V_0)^{1/\pi}$$

$$= -I - \ln\left(\frac{V (1 + \pi V_0)^{1/\pi}}{V_0^{1/(1 + \pi V_0)}}\right)$$

$$\Rightarrow \frac{V (1 + \pi V_0)^{1/\pi}}{V_0^{1/(1 + \pi V_0)}} = \exp(-X_0 \tau) \exp(-I)$$

$$\Rightarrow V(\tau) \approx V_0 \frac{1}{1 + \pi V_0} \frac{1}{(1 + \pi V_0)^{1/\pi}} \exp\left(-\frac{X_0}{(1 + \pi V_0)^{1/\pi}} \tau\right) \exp\left(-I / (1 + \pi V_0)^{1/\pi}\right)$$

comme $\frac{1}{\pi} + 1 \sim 5$, alors
 $1/(1/\pi + V)^{1/\pi+1} \sim 1/(1/\pi)^{1/\pi+1}$ si $\tau \rightarrow \infty$
 (c.e. impact pour $V=0$): ce terme n'engendre pas de constante du type $e^{-\dots}$ mais contribue à la pente $e^{-\dots}$ qui dans cette approximation a été vérifiée numériquement.

$$\Rightarrow \boxed{V(\tau) \approx V_0 \frac{1/(1+V_0 e)^{1/(1+k)}}{e^{-\frac{x_0}{x_0} I} e^{-x_0 \tau}} ; I = \int_0^{V_0} ds h(s) \frac{d^2}{ds^2} \left[-\left(\frac{1/(1+k) + V_0}{1/(1+k) + s} \right)^{1/(1+k)} \right]} \quad (11)$$

$$k = 4/\pi - 1 = 0.2732\dots$$

$$x_\infty = \frac{x_0}{(1+V_0 e)^{1/(1+k)}}$$

Comportement asympt. de X et Y:

$$(9) \Rightarrow \frac{d}{d\tau} \left(\frac{1}{X} \right) \approx f \exp(-x_0 \tau) ; \boxed{f = V_0 \frac{1/(1+V_0 e)^{1/(1+k)}}{e^{-\frac{x_0}{x_0} I}} ; \int_2^\infty}$$

$$\Rightarrow \frac{1}{x_\infty} - \frac{1}{X(\tau)} \approx f e^{-x_0 \tau} \Big|_2^\infty = \frac{f}{x_\infty} e^{-x_0 \tau}$$

$$\Rightarrow \frac{1}{X(\tau)} \approx \frac{1 - f e^{-x_0 \tau}}{x_\infty}$$

$$\Rightarrow \boxed{X(\tau) \approx \frac{x_\infty}{1 - f e^{-x_0 \tau}} \approx x_\infty \left[1 + f e^{-x_0 \tau} \right]} \quad (12)$$

et:

$$(5) \Rightarrow \dot{X} = -XY, \text{ avec } \dot{X} = x_\infty \left[0 - x_\infty f e^{-x_0 \tau} \right]$$

$$\Rightarrow -x_\infty f e^{-x_0 \tau} \approx -x_\infty \left[1 + f e^{-x_0 \tau} \right] Y$$

$$\Rightarrow Y(\tau) \approx \frac{x_\infty f e^{-x_0 \tau}}{1 + f e^{-x_0 \tau}} \approx x_\infty f e^{-x_0 \tau} \left[1 - f e^{-x_0 \tau} \right]$$

$$\Rightarrow \boxed{Y(\tau) \approx x_\infty f e^{-x_0 \tau}} \quad (13)$$

Relation entre X et Y:

$$\boxed{X(\tau) = x_\infty + Y(\tau)} \quad (14)$$

Borne par Y: (32) P1a. \Rightarrow

$$\frac{dx}{dy} = \frac{4\delta x^2 + (3+\delta^2)xy}{4y^2 + (3+\delta^2)xy}$$

$$; Y = v \cdot x \Rightarrow x = \frac{Y}{v}$$

$$\frac{dx}{dy} = \frac{d}{dy}\left(\frac{Y}{v}\right) = \frac{1}{v} - Y \frac{1}{v^2} \frac{dv}{dy}$$

$$\Rightarrow \left(1 - \frac{Y}{v} \frac{dv}{dy}\right) = v \frac{4\delta \left(\frac{Y}{v}\right)^2 + (3+\delta^2)Y\left(\frac{Y}{v}\right)}{4Y^2 + (3+\delta^2)Y\left(\frac{Y}{v}\right)}$$

$$= \frac{4\delta + (3+\delta^2)v}{4v + (3+\delta^2)}$$

$$\Rightarrow -\frac{Y}{v} \frac{dv}{dy} = \frac{4\delta + (3+\delta^2)v - 4v - (3+\delta^2)}{4v + (3+\delta^2)} = \frac{4\delta - 3 - \delta^2 + v[3 + \delta^2 - 4]}{4v + (3+\delta^2)}$$

$$\Rightarrow \frac{dy}{Y} = \frac{4v + (3+\delta^2)}{[3 + \delta^2 - 4\delta + v(1 - \delta^2)]v} dv = \left(\frac{\alpha}{v} + \frac{\beta(1 - \delta^2)}{(1 - \delta^2)v + 3 + \delta^2 - 4\delta}\right) dv ; \alpha, \beta \text{ à déterminer (1)}$$

$$\underline{\alpha, \beta}: \alpha \cdot ((1 - \delta^2)v + 3 + \delta^2 - 4\delta) + \beta(1 - \delta^2)v = 4v + (3 + \delta^2)$$

$$\Rightarrow \begin{cases} \alpha(1 - \delta^2) + \beta(1 - \delta^2) = 4 \\ \alpha(3 + \delta^2 - 4\delta) = 3 + \delta^2 \end{cases} \Rightarrow \alpha = \frac{3 + \delta^2}{3 + \delta^2 - 4\delta} = \frac{3 + \delta^2}{(1 - \delta)(3 - \delta)} \quad (2)$$

$$\Rightarrow \beta = \frac{4 - \alpha(1 - \delta^2)}{1 - \delta^2} = \frac{4 - (3 + \delta^2)(1 - \delta^2) / ((1 - \delta)(3 - \delta))}{(1 - \delta^2)}$$

$$= \frac{4(3 - \delta) - (3 + \delta^2)(1 + \delta)}{(1 - \delta^2)(3 - \delta)} \quad (3)$$

(2) et (3) dans (1) \Rightarrow

$$\frac{dy}{Y} = \left[\frac{\alpha}{v} + \frac{\beta(1 - \delta^2)}{(1 - \delta^2)v + 3 + \delta^2 - 4\delta} \right] dv ; \alpha = \frac{3 + \delta^2}{(1 - \delta)(3 - \delta)} ; \beta = \frac{4(3 - \delta) - (3 + \delta^2)(1 + \delta)}{(1 - \delta^2)(3 - \delta)} \quad (4)$$

Intégration: tot. t:

$$\int_{\frac{y_0}{Y}}^{\frac{y_0}{Y}} \frac{dz}{z} = \int_{\frac{v_0}{v}}^{\frac{v_0}{v}} dz \frac{\alpha}{z} + \int_{\frac{v_0}{v}}^{\frac{v_0}{v}} dz \frac{\beta(1 - \delta^2)}{(1 - \delta^2)z + (1 - \delta)(3 - \delta)}$$

$$= \ln\left(\frac{y_0}{Y}\right) = \alpha \ln\left(\frac{v_0}{v}\right) = \int_{\frac{v_0}{v}}^{\frac{v_0}{v}} dz \frac{\beta}{z + \frac{(1 - \delta)(3 - \delta)}{(1 - \delta^2)(1 + \delta)}} = \beta \ln\left(\frac{v_0 + \frac{3 - \delta}{1 + \delta}}{v + \frac{3 - \delta}{1 + \delta}}\right)$$

$$\Rightarrow \frac{y_0}{Y} = \left(\frac{v_0}{v}\right)^\alpha \left(\frac{v_0 + \frac{3 - \delta}{1 + \delta}}{v + \frac{3 - \delta}{1 + \delta}}\right)^\beta \quad (5)$$

Avec:

$$\frac{d}{dt}\left(\frac{1}{Y}\right) = 4 + (3 + \delta^2) \frac{x}{Y}$$

(5) devient:

$$\frac{d}{dt}\left(\frac{y_0}{Y}\right) = y_0 \left[4 + (3 + \delta^2) \frac{1}{v} \right] = \frac{d}{dt} \left[\left(\frac{v_0}{v}\right)^\alpha \left(\frac{v_0 + \frac{3 - \delta}{1 + \delta}}{v + \frac{3 - \delta}{1 + \delta}}\right)^\beta \right]$$

$$= \frac{d}{dv} \left[\left(\frac{v_0}{v}\right)^\alpha \left(\frac{v_0 + \frac{3 - \delta}{1 + \delta}}{v + \frac{3 - \delta}{1 + \delta}}\right)^\beta \right] \frac{dv}{dt}$$

$$\Rightarrow y_0 \tau = \frac{dv}{4 + \frac{3 + \delta^2}{v}} \frac{d}{dv} \left\{ \left(\frac{v_0}{v}\right)^\alpha \left(\frac{v_0 + \frac{3 - \delta}{1 + \delta}}{v + \frac{3 - \delta}{1 + \delta}}\right)^\beta \right\}$$

$$\Rightarrow y_0 \tau = \int_{v_0}^v \frac{du}{4 + \frac{3 + \delta^2}{u}} \frac{d}{du} \left\{ \left(\frac{v_0}{u}\right)^\alpha \left(\frac{v_0 + \frac{3 - \delta}{1 + \delta}}{u + \frac{3 - \delta}{1 + \delta}}\right)^\beta \right\} ; \alpha > 0, \beta > 0$$

$$\Rightarrow \boxed{y_0 \tau = \int_{v_0}^v du \frac{u}{3 + \delta^2 + 4u} \left\{ -\frac{d}{du} \left[\left(\frac{v_0}{u}\right)^\alpha \left(\frac{v_0 + \frac{3 - \delta}{1 + \delta}}{u + \frac{3 - \delta}{1 + \delta}}\right)^\beta \right] \right\}} \quad (6)$$

Barre supérieure:

$$\frac{u}{3+\delta^2+4u} \leq \frac{v_0}{3+\delta^2+4v} \leq \frac{v_0}{3+\delta^2}$$

$$\Rightarrow y_0 \tau \leq \frac{v_0}{3+\delta^2} \left(-1 + \frac{y_0}{\tau} \right)$$

$$\Rightarrow Y(\tau) \leq \frac{y_0/x_0 \cdot y_0}{\frac{y_0}{x_0} + (3+\delta^2)Y_0\tau} = \frac{y_0}{1 + (3+\delta^2)x_0\tau}$$

$$\begin{cases} \dot{x} = -4\delta x^2 - (3+\delta^2)xy \\ \dot{y} = -4y^2 - (3+\delta^2)xy \end{cases} \rightarrow \left(\frac{y}{x}\right) = 4\delta + (3+\delta^2)\frac{y}{x} = 4\delta + (3+\delta^2)v \quad (*)$$

$$\Rightarrow \frac{dy}{dx} = \frac{4y^2 + (3+\delta^2)xy}{4\delta x^2 + (3+\delta^2)xy} \quad (1) ; Y = v \cdot X \quad (2) \text{ (def. de } v)$$

$$\Rightarrow \frac{dy}{dx} \stackrel{(2)}{=} v + \frac{dv}{dx} \stackrel{(1)}{=} \frac{4v^2 + (3+\delta^2)v}{4\delta + (3+\delta^2)v}$$

$$\Rightarrow x \frac{dv}{dx} = \frac{4v^2 + (3+\delta^2)v - 4\delta v - (3+\delta^2)v^2}{4\delta + (3+\delta^2)v} = \frac{v^2 - \delta^2 v^2 + (3+\delta^2 - 4\delta)v}{4\delta + (3+\delta^2)v} = \frac{(1-\delta^2)v^2 + (3+\delta^2 - 4\delta)v}{4\delta + (3+\delta^2)v}$$

$$\Rightarrow \frac{dx}{x} = \frac{4\delta + (3+\delta^2)v \, dv}{[(1-\delta^2)v^2 + (3+\delta^2 - 4\delta)]v} ; 3+\delta^2 - 4\delta = (1-\delta)(3-\delta)$$

$$\Rightarrow (1-\delta) \frac{dx}{x} = \frac{4\delta + (3+\delta^2)v}{[(1+\delta)v + (3-\delta)]v} \frac{dv}{v} = \frac{\alpha}{v} + \frac{\beta}{(1+\delta)v + (3-\delta)}$$

$$\Rightarrow \frac{\alpha((1+\delta)v + (3-\delta)) + \beta v}{[(1+\delta)v + (3-\delta)]v} = \frac{4\delta + (3+\delta^2)v}{[(1+\delta)v + (3-\delta)]v}$$

$$\Rightarrow \begin{cases} \alpha(3-\delta) = 4\delta \\ \beta + \alpha(1+\delta) = (3+\delta^2) \end{cases}$$

$$\Rightarrow \begin{cases} \alpha = \frac{4\delta}{3-\delta} \\ \beta = 3+\delta^2 - \frac{4\delta(1+\delta)}{3-\delta} \end{cases}$$

Equation suite: $(1-\delta) \ln\left(\frac{x}{x_0}\right) = \alpha \ln\left(\frac{v}{v_0}\right) + \frac{\beta}{1+\delta} \ln\left(\frac{(1+\delta)v + 3-\delta}{(1+\delta)v_0 + 3-\delta}\right)$

$$; v_0 = \frac{y_0}{x_0}$$

$$\Rightarrow \left(\frac{x}{x_0}\right)^{1-\delta} = \left(\frac{v}{v_0}\right)^\alpha \left(\frac{(1+\delta)v + 3-\delta}{(1+\delta)v_0 + 3-\delta}\right)^{\frac{\beta}{1+\delta}}$$

$$\Rightarrow \frac{x_0}{x} = \left(\frac{v}{v_0}\right)^{-\frac{\alpha}{1-\delta}} \left(\frac{(1+\delta)v + 3-\delta}{(1+\delta)v_0 + 3-\delta}\right)^{-\frac{\beta}{1+\delta}}$$

$$\Rightarrow x_0 \left(\frac{1}{x}\right)^{\frac{1}{1-\delta}} = \frac{d}{dz} \left[\underbrace{\left(\frac{v}{v_0}\right)^{-\frac{\alpha}{1-\delta}} \left(\frac{(1+\delta)v + 3-\delta}{(1+\delta)v_0 + 3-\delta}\right)^{-\frac{\beta}{1+\delta}}}_{= \chi(v)} \right] = \frac{d}{dz} \chi(v)$$

$$\Rightarrow x_0 (4\delta + (3+\delta^2)v) = \frac{d}{dz} \chi(v)$$

$$\Rightarrow dz = \frac{\chi(v)}{x_0(4\delta + (3+\delta^2)v)}$$

$$\Rightarrow \boxed{\chi = \int_{v_0}^v \frac{\chi(u)}{x_0[4\delta + (3+\delta^2)u]} du}$$

num \rightarrow num \rightarrow num
 $\tau = \tau(v) \rightarrow v = v(\tau) \rightarrow \frac{Y(\tau)}{X(\tau)} = \text{denom} + \text{intégrer}$
 pris dans (*)
 + comportement asymptotique : terme divergeant dominant.

$$\mathcal{Z} = \int_{v_0}^v du \frac{\chi(u)}{x_0 [4\delta(1+\delta^2)u]} = \frac{1}{x_0} \int_{v_0}^v du \left(\frac{u}{v_0}\right)^{\frac{1-\delta}{\alpha}} \left[\frac{(1+\delta)u + 3-\delta}{(1+\delta)v_0 + 3-\delta} \right]^{\frac{1-\delta}{\beta}} \frac{1}{4\delta + (3+\delta)u}$$

avec: $\alpha = \frac{4\delta}{3-\delta}$; $\alpha \in]0, 2[$
 $\beta = 3+\delta^2 - \frac{4\delta(1+\delta)}{3-\delta} \in]0, 3[$

Calcul: exposants non entiers avec fractions \Rightarrow fct. hypergéométriques... $\Rightarrow \mathcal{E} \Rightarrow$ recurs logiciel:

$$\begin{aligned} \Rightarrow \mathcal{Z}(v) &= \frac{1}{4\delta(\alpha+\delta-1)x_0} \left[(\delta-1) \left(\frac{3+\delta(-1+v)+v}{3-\delta}\right)^{\frac{\beta}{\delta-1}} \left(\frac{3+\delta(-1+v_0)+v_0}{3-\delta}\right)^{\frac{\beta}{\delta-1}} \cdot \right. \\ &\quad \left. v \left(\frac{v}{v_0}\right)^{\frac{\alpha}{-1+\delta}} \left(\frac{3+\delta(-1+v)+v}{3+\delta(-1+v_0)+v_0}\right)^{-\frac{\beta}{-1+\delta^2}} \left(\frac{3+\delta(-1+v_0)+v_0}{3-\delta}\right)^{\frac{\beta}{-1+\delta^2}} \text{Appel } F_1[\dots] \right. \\ &\quad \left. - \left(\frac{3+\delta(-1+v)+v}{3-\delta}\right)^{-\frac{\beta}{-1+\delta^2}} v_0 \text{Appel } F_1[\dots] \right] \\ &= \frac{1}{4\delta(\alpha+\delta-1)x_0} \left[(-1+\delta) \left\{ v \left(\frac{v}{v_0}\right)^{\frac{\alpha}{-1+\delta}} \left(1 - \frac{(1+\delta)v}{-3+\delta}\right)^{-\frac{\beta}{-1+\delta^2}} \left(\frac{3+\delta(-1+v)+v}{3+\delta(-1+v_0)+v_0}\right)^{-\frac{\beta}{-1+\delta^2}} \text{Appel } F_1(\dots) \right. \right. \\ &\quad \left. \left. + v_0 \left(1 - \frac{(1+\delta)v_0}{-3+\delta}\right)^{-\frac{\beta}{-1+\delta^2}} \text{Appel } F_1(-v_0) \right\} \right] \\ &= \frac{\delta-1}{4\delta(\alpha+\delta-1)x_0} \left[v \left(\frac{v}{v_0}\right)^{\frac{\alpha}{-1+\delta}} \left(1 + \frac{(1+\delta)v}{3-\delta}\right)^{\frac{\beta}{-1+\delta^2}} \left(\frac{3-\delta+v(\delta+1)}{3-\delta+v_0(\delta+1)}\right)^{-\frac{\beta}{-1+\delta^2}} \text{App.} \right. \\ &\quad \left. + v_0 \left(1 + \frac{(1+\delta)v_0}{3-\delta}\right)^{\frac{\beta}{-1+\delta^2}} \text{App.} \right] \\ &= \frac{\delta-1}{4\delta(\alpha+\delta-1)x_0} \left[v \left(\frac{v}{v_0}\right)^{\frac{\alpha}{\delta-1}} \left(\frac{3-\delta-(1+\delta)v}{3-\delta} \cdot \frac{3-\delta+v_0(\delta+1)}{3-\delta+(1+\delta)v}\right)^{\frac{\beta}{-1+\delta^2}} \text{App.} \right. \\ &\quad \left. + v_0 \cdot \left(\frac{3-\delta+(1+\delta)v_0}{3-\delta}\right)^{\frac{\beta}{-1+\delta^2}} \text{App} \right] \end{aligned}$$

~~$\frac{1-\delta}{x_0 4\delta(1-\delta-\alpha)} v \left(\frac{v}{v_0}\right)^{\frac{\alpha}{\delta-1}} \left(1 - \frac{1+\delta}{3-\delta} v\right)$~~

~~$\mathcal{Z}(v) = \frac{1-\delta}{4\delta x_0(1-\delta-\alpha)}$~~

~~$\mathcal{Z}(v) = \frac{\delta-1}{4\delta(\alpha+\delta-1)x_0} v \left(\frac{v}{v_0}\right)^{\frac{\alpha}{-1+\delta}}$~~

$$\mathcal{Z}(v) = \frac{1-\delta}{x_0 4\delta(1-\delta-\alpha)} \left[v \left(\frac{v}{v_0}\right)^{\frac{\alpha}{-1+\delta}} \left(\frac{3-\delta-(1+\delta)v}{3-\delta+(1+\delta)v} \cdot \frac{3-\delta+v_0(\delta+1)}{3-\delta}\right)^{\frac{\beta}{-1+\delta^2}} \text{Appel } F_1 \left[1 - \frac{\alpha}{-1+\delta}, \frac{\beta}{-1+\delta^2}, 1, 2 - \frac{\alpha}{-1+\delta}, \frac{(1+\delta)v}{\delta-3}, -\frac{(3+\delta^2)v}{4\delta} \right] \right. \\ \left. + v_0 \left(\frac{3-\delta+v_0(1+\delta)}{3-\delta}\right)^{\frac{\beta}{-1+\delta^2}} \text{Appel } F_1[\dots, v_0, v_0] \right]$$

Soit: $h(\delta, v_0) = \left[\frac{3-\delta + v_0(\delta+1)}{3-\delta} \right]^{\frac{\beta}{1-\delta^2}}$

$App[V] = App F_1 \left[\left(\frac{1-\frac{\alpha}{1-\delta}}{\alpha_1} \right), \frac{\beta}{1-\delta^2}, \left(\frac{1}{1+\alpha_1} \right), \left(\frac{2-\frac{\alpha}{1-\delta}}{1+\alpha_1} \right), -\frac{(1+\delta)}{3-\delta} V, -\frac{3+\delta^2}{4\delta} V \right]$ (Gradshteyn p. 1080 ss.)

Alors:

$$\tau(v) = \frac{1-\delta}{4\delta(1-\delta-\alpha)} \frac{1}{x_0} \cdot \left\{ v \left(\frac{v_0}{v} \right)^{\frac{\alpha}{1-\delta}} \left(\frac{3-\delta-(1+\delta)v}{3-\delta+(1+\delta)v} \right)^{\frac{\beta}{1-\delta^2}} \cdot h(\delta, v_0) App[V] + v_0 h(\delta, v_0) App[v_0] \right\}$$

$$\Rightarrow \tau(v) = \frac{1}{x_0} \cdot \frac{h(\delta, v_0) \cdot (1-\delta)}{4\delta(1-\delta-\alpha)} \left\{ v \left(\frac{v_0}{v} \right)^{\frac{\alpha}{1-\delta}} \left(\frac{3-\delta-(1+\delta)v}{3-\delta+(1+\delta)v} \right)^{\frac{\beta}{1-\delta^2}} App[V] + v_0 App[v_0] \right\}$$

avec: $App F_1[a; b_1, b_2; c; x, y]$ qui est la fonction hypergéométrique de deux variables (Gradshteyn p. 1080 ss.)